

Interaction Notes

Note 240

February 1975

Penetration of Electromagnetic Pulses Through
Larger Apertures in Shielded Enclosures

R. Mittra
L. Wilson Pearson
Electromagnetics Laboratory
Department of Electrical Engineering
University of Illinois at Urbana-Champaign
Urbana, Illinois 61801

Abstract

The results of an initial investigation of the Singularity Expansion Method representation of the electromagnetic coupling through a rectangular aperture in a perfectly conducting sheet are reported. The problem is formulated in terms of the coupled Hallén-type integral equations for the dual problem of a rectangular plate. The integral equations are converted to a system of linear algebraic equations by way of the method of moments with subsectionally constant expansion functions and collocation testing. Several techniques used in minimizing the execution time of the computations are described. Some difficulties in accurately approximating the singularities of the system of integral equations by the singularities of the algebraic system are discussed. These difficulties arise because the subsectionally constant representation for the current cannot adequately represent the correct edge singularities in the currents on the plate. A set of pole trajectories indicative of the trends in pole location for the plate is reported. A listing of the pertinent computer code is provided.

This study was performed under subcontract to

The Dikewood Corporation
1009 Bradbury Drive, S.E.
University Research Park
Albuquerque, New Mexico 87106

and has been fully supported by the Defense Nuclear Agency (DNA) under

DNA Subtask EB088
EMP Interaction and Coupling
DNA Work Unit 33
Coupling Characteristics of Apertures

CONTENTS

<u>Section</u>		<u>Page</u>
I	INTRODUCTION	5
II	THIN-PLATE INTEGRAL EQUATION FORMULATION FOR COMPLEX WAVENUMBER	7
III	SYMMETRY CONDITIONS FOR THE NATURAL MODE CURRENTS	10
IV	THE NUMERICAL MODEL	16
V	ALGORITHMIC CONSIDERATIONS IN EVALUATING THE SYSTEM DETERMINANT	21
VI	NUMERICAL CHECKS ON THE ACCURACY OF THE POLES	25
VII	POLE TRAJECTORIES AS A FUNCTION OF SHAPE RATIO	31
VIII	CONCLUSIONS	33
	APPENDIX A: THE SELF-PATCH INTEGRATION	35
	APPENDIX B: THE SPARSE MATRIX ALGORITHMS	37
	APPENDIX C: PROGRAM LISTINGS	43
	REFERENCES	74

ILLUSTRATIONS

<u>Figure</u>		<u>Page</u>
1	Geometry of the Rectangular Plate	8
2	Lowest Order Natural Mode Current Pairs for Each of the Symmetry Cases, a) J_x Symmetric w.r.t. x-Axis and Symmetric w.r.t. y-Axis, b) Symmetric-Antisymmetric, c) Antisymmetric-Symmetric, and d) Antisymmetric-Antisymmetric	15
3	Subsectioning for the Discretization of the Integral Equations	17
4	Organization of the System of Linear Equations	19
5	a) Conceptual Zoning for Calculation of the Interaction Matrix, b) Example of the Four Interaction Contributions to a Single Source Term	22
6	Calculated Pole Locations for Thin-Strip for Varying Numbers of Zones in the x-Direction and Different Edge Treatments (Cylinder Results from Ref. 6)	26
7	Computer Pole Trajectory Under Varying w/L with Zoning Changes	28
8	Behavior of Singular Component of Current at the Edge Under Change in Transverse Zoning	30
9	Pole Trajectories as Computed with 4 x 3 Zoning	32
B1	Submatrix Organization for the Sparse Matrix Algorithms, a) the Original Matrix, and b) Triangularized Form with the Generated CMAT4	38

TABLES

<u>Table</u>		<u>Page</u>
1	Compatible Current Symmetry Features	14
2	Matrix Formulation Parameters	20
B1	Primary Indexing Quantities in the Algorithm	42

SECTION I
INTRODUCTION

This report presents the results of an investigation for representing the transient electromagnetic coupling through a rectangular aperture by means of the singularity expansion method.

The singularity expansion method, which was introduced by Baum in 1971 (ref. 1), has been subsequently applied to many scatterer geometries. The essence of the singularity expansion method is the representing of the temporal response of a body in terms of the complex natural frequencies for the body.

Taylor et al. point out that the singularity expansion for an aperture in an infinite perfectly conducting screen can be determined in terms of that for the complementary perfectly conducting plate by way of Babinet's principle (ref. 2). This approach was taken in the work reported here. The remaining discussion is directed to the equivalent problem of determining the current distributions on the complementary plate geometry.

Rahmat-Samii and Mittra have derived a coupled pair of Hallen-type integral equations governing the current behavior on the rectangular plate (ref. 3). The work reported here builds on their work by generalizing the integral equations and solution method to the complex frequency plane for the

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1. Baum, C. E., "On the Singularity Expansion Method for the Solution of Electromagnetic Interaction Problems," Interaction Note 88, Air Force Weapons Laboratory, Kirtland AFB, NM, December 1971.
 2. Taylor, C. D., Crow, T. T., and Chen, K-T, "On the Singularity Expansion Method Applied to Aperture Penetration: Part I Theory," Interaction Note 134, Air Force Weapons Laboratory, Kirtland AFB, NM, May 1973.
 3. Rahmat-Samii, Y. and Mittra, R., "Integral Equation Solution and RCS Computation of a Thin Rectangular Plate," Interaction Note 156, Air Force Weapons Laboratory, Kirtland AFB, NM, December 1973.

SEM application. The same method-of-moments formulation, as described in (ref. 3), is used, i.e., two-dimensional pulse expansion functions with collocation testing.

In order that the computation time be practical in a problem of this complexity, a great deal of care was given to algorithmic streamlining in the conduct of this work. The streamlining includes maximum exploitation of geometric symmetry, organization of calculations to make use of redundant terms and partial terms occurring in the calculation, and direct algorithmic exploitation of matrix **sparseness**. The end result is a highly efficient computer code. Key features of the algorithms are discussed in this report.

The pulse expansion appears to be inadequate in accurately modeling the rectangular plate. The difficulty, which relates to representing the actual size of the plate, is demonstrated and discussed herein. Remedies for the problem are suggested, but they are outside the scope of the present work.

By holding the zoning scheme invariant while the aspect ratio of the plate was changed, self-consistent pole trajectories for the four fundamental modes were determined. For the reasons cited above, these poles cannot claim to be the exact poles for the body. They are, however, indicative of the trends in the pole behavior for the plate under change in aspect ratio. These results are reported and discussed in this context.

SECTION II

THIN-PLATE INTEGRAL EQUATION FORMULATION FOR COMPLEX WAVENUMBER

Rahmat-Samii and Mittra (ref. 3) give an integral equation formulation for the rectangular plate subject to time-harmonic excitation. Their results may be directly extended to the complex wavenumber case. That is, for the geometry in Figure 1 with $\exp[st]$ time dependence, $s = \sigma + j\omega$ complex, and an incident plane-wave magnetic field component

$\vec{H}^i = [H_{ox}^i \hat{u}_x + H_{oy}^i \hat{u}_y + H_{oz}^i \hat{u}_z] \exp[j(k_x x + k_y y + k_z z)]$ the following coupled integral equations result:

$$\int_{-L/2}^{L/2} \int_{-w/2}^{w/2} \begin{Bmatrix} J_x(x,y) \\ J_y(x,y) \end{Bmatrix} K(x,y|x',y') dx' dy' = \frac{j}{k_z} \begin{Bmatrix} H_{og}^i \\ -H_{ox}^i \end{Bmatrix} \exp[j(k_x x + k_y y)]$$

$$+ \frac{\pi}{k} \begin{Bmatrix} -1 \\ -j \end{Bmatrix} \sum_{-\infty}^{\infty} C_n [j^{n+1} \exp[j(n+1)\phi] J_{n+1}(-s\rho/c)$$

$$+ \begin{Bmatrix} 1 \\ -1 \end{Bmatrix} j^{n-1} \exp[j(n-1)\phi] J_{n-1}(-s\rho/c)] \quad (1)$$

The kernel is given by

$$K(x,y | x',y') = \exp[-sR/c]/R \quad (2)$$

with

$$R = [(x - x')^2 + (y - y')^2]^{1/2}$$

The $J_x(x,y)$ and $J_y(x,y)$ denote the respective x and y components of current on the plate; $J_n(\zeta)$ denotes the Bessel function of the first kind; the C_n are unknown constants; c is the velocity of light; and (ρ, ϕ) are the polar coordinates for the point (x,y) on the plate. Equation (1) holds for $x \in (-L/2, L/2)$ and $y \in (-w/2, w/2)$, and $z = 0$.

It is pointed out that the two integral equations represented by (1) are

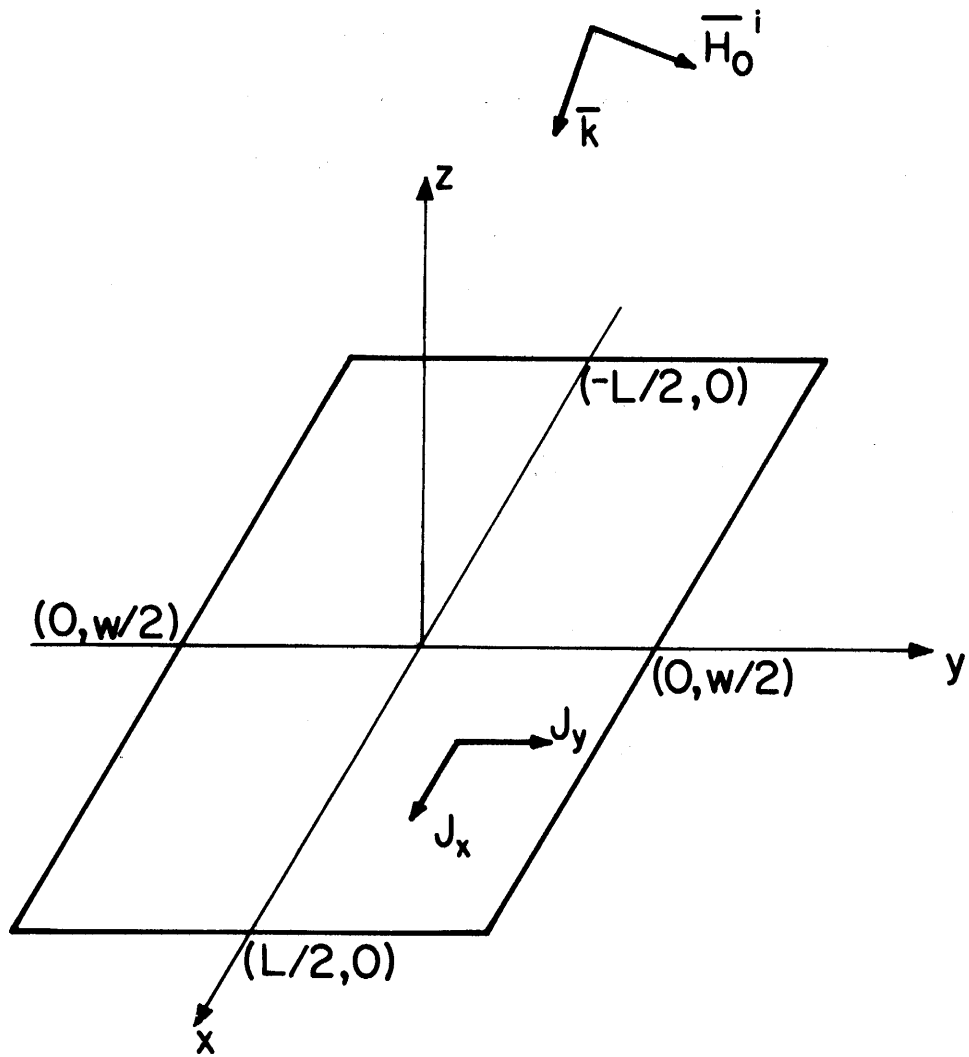


Figure 1. Geometry of the Rectangular Plate

coupled through the C_n in the summation in the right-hand side. This summation is simply a Bessel function expansion of the homogeneous solution to the wave equation which occurs in the derivation of (1). Details of arriving at this expansion are found in (ref. 3). The pair of integral equations (1) is complete in the sense that no additional constraints are needed to correctly specify the currents. It is noteworthy, however, that current solutions to (1) satisfy the Meixner's edge condition (ref. 4). Namely,

$$\left. \begin{aligned}
 J_x[\pm(L/2 - d), y] &\rightarrow d^{1/2} \\
 J_y[\pm(L/2 - d), y] &\rightarrow d^{-1/2} \\
 J_x[x, \pm(w/2 - d)] &\rightarrow d^{-1/2} \\
 J_y[x, \pm(w/2 - d)] &\rightarrow d^{1/2}
 \end{aligned} \right\} \quad d \rightarrow 0 \quad (3)$$

The ability of a numerical solution to approximate the behavior of eqn. (3) is a key point in a subsequent discussion.

4. Bladel, J. Van, Electromagnetic Fields, McGraw-Hill, New York, pp. 385-387, 1964.

SECTION III

SYMMETRY CONDITIONS FOR THE NATURAL MODE CURRENTS

The natural frequencies of (1) occur when the complex frequency s is such that there are non-trivial J_x and J_y and the accompanying C_n satisfying (1) for $\bar{H}^i = 0$. Such J_x and J_y solutions are natural mode current solutions for the rectangular plate, and the concomitant value of s is a pole of the plate. The vanishing of incident wave dependence gives rise to symmetry in the integral equations. By discerning the symmetry relations a priori and bringing them to bear upon solution procedures, one gains significant computational savings in the numerical solution for poles and natural modes. These symmetry relations are explored in this section.

The excitation-free form of (1) is

$$\int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_x K(x,y|x',y') dx' dy' = \frac{j\pi c}{s} \sum_{-\infty}^{\infty} C_n \left\{ j^{n+1} \exp[j(n+1)\phi] J_{n+1}(-s\rho/c) + j^{n-1} \exp[j(n-1)\phi] J_{n-1}(-s\rho/c) \right\} \quad (4a)$$

and

$$\int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_y K(x,y|x',y') dx' dy' = \frac{\pi c}{s} \sum_{-\infty}^{\infty} C_n \left\{ j^{n+1} \exp[j(n+1)\phi] J_{n+1}(-s\rho/c) - j^{n-1} \exp[j(n-1)\phi] J_{n-1}(-s\rho/c) \right\} \quad (4b)$$

By using the symmetry of the Bessel function with respect to order, expanding the exponentials by way of Euler's identity, and appropriately adjusting the indices, one arrives at the following equation after some manipulation.

$$\begin{aligned}
& \int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_x K \, dx' \, dy' \\
&= \frac{j\pi c}{s} \sum_{n=0}^{\infty} \left\{ j^{n+1} d_n^+ [\cos(n+1)\phi J_{n-1}(-s\rho/c) - u_{n-1} \cos(n-1)\phi J_{n-1}(-s\rho/c)] \right. \\
&\quad \left. - j^n d_n^- [\sin(n+1)\phi J_{n+1}(-s\rho/c) - \sin(n-1)\phi J_{n-1}(-s\rho/c)] \right\} \quad (5a)
\end{aligned}$$

and

$$\begin{aligned}
& \int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_y K \, dx' \, dy' \\
&= \frac{j\pi c}{s} \sum_{n=0}^{\infty} \left\{ j^{n+1} d_n^+ [\sin(n+1)\phi J_{n+1}(-s\rho/c) + \sin(n-1)\phi J_{n-1}(-s\rho/c)] \right. \\
&\quad \left. + j^n d_n^- [\cos(n+1)\phi J_{n+1}(-s\rho/c) + u_{n-1} \cos(n-1)\phi J_{n-1}(-s\rho/c)] \right\} \quad (5b)
\end{aligned}$$

where

$$d_n^{\pm} = C_n \pm C_{-n}$$

and

$$u_n = \begin{cases} 1, & n \geq 0 \\ 0, & n < 0 \end{cases}$$

It is noted that the d_n^+ multiply terms containing cosine functions in the J_x equation, while they multiply terms containing sine functions in the J_y equation. The situation is reversed for the d_n^- .

Because of the symmetry properties of the kernel, the integral operator on the left-hand sides of (5) produces a function whose symmetry character is identical to that of the current on which it operates. Then, for a given current symmetry, only part of the d_n^{\pm} on the right-hand side may be non-zero because of the symmetries possessed by the trigonometric terms. Thus, the respective symmetries for J_x and J_y , which are compatible, and the

surviving terms in the right-side series may be discerned by 1) postulating a symmetry for J_x , 2) determining from (5a) which right-hand side terms survive so as to be compatible with the J_x symmetry, 3) observing in (5b) the variation which terms have non-zero coefficients, and 4) determining the J_y symmetry conditions compatible with the postulated J_x symmetry conditions.

For example, if J_x is symmetric with respect to the y axis and anti-symmetric with respect to the x axis, only $\sin(n+1)\phi$ terms with n even are compatible in (5a). Thus, only d_n^- , n even, may be non-zero. In the right-hand side of (5b), the coefficients multiply $\cos(n+1)\phi$ terms with n even. These cosines sum to functions which are antisymmetric with respect to the y axis and symmetric with respect to the x axis. Stated mathematically, if

$$J_x(x,y) = J_x(-x,y) \quad (6a)$$

and

$$J_x(x,y) = -J_x(x,-y) \quad (6b)$$

then

$$d_n^+ = 0, \quad \text{for all } n, \quad (6c)$$

$$d_n^- = 0, \quad n \text{ odd}, \quad (6d)$$

and

$$J_y(x,y) = -J_y(-x,y) \quad (6e)$$

$$J_y(x,y) = J_y(x,-y) \quad (6f)$$

These vector symmetries are in accord with the general symmetry relations given by Baum (ref. 5). The information in (6) may be used to reduce the complexity of the integral equations (4), viz., by (6a,b,e,f) the range of each integration may be halved while by (6c,d) the zero terms of the right-hand side are known a priori:

$$\begin{aligned} & \int_0^{L/2} \int_0^{w/2} J_x K^{-+}(x,y|x',y') dx' dy' \\ &= \frac{\pi c}{s} \sum_{\substack{n=0 \\ n \text{ even}}}^{\infty} d_n^- j^{n-1} [\sin(n+1)\phi J_{n-1}(-s\rho/c) - \sin(n-1)\phi J_{n-1}(-s\rho/c)] \end{aligned} \quad (7a)$$

5. Baum, C. E., "Interaction of Electromagnetic Fields with any Object which has an Electromagnetic Symmetry Plane," Interaction Note 63, Air Force Weapons Laboratory, Kirtland AFB, NM, March 1971.

and

$$\int_0^{L/2} \int_0^{w/2} J_y K^{\pm}(x,y|x',y') dx' dy'$$

$$= \frac{\pi c}{s} \sum_{\substack{n=0 \\ n \text{ even}}}^{\infty} j^{n+1} d_x^- [\cos(n+1)\phi J_{n+1}(-sp/c) + u_{n-1} \cos(n-1)\phi J_{n-1}(-sp/c)] \quad (7b)$$

where

$$K^{\pm}(x,y|x',y') = K(x,y|x',y') - K(x,y|-x',y') \\ + K(x,y|x',-y') - K(x,y|-x',-y') \quad (8a)$$

and

$$K^{\mp}(x,y|x',y') = K(x,y|x',y') + K(x,y|-x',y') \\ - K(x,y|x',-y') - K(x,y|-x',-y') \quad (8b)$$

For subsequent reference

$$K^{++}(x,y|x',y') = K(x,y|x',y') + K(x,y|-x',y') \\ + K(x,y|x',-y') + K(x,y|-x',-y') \quad (8c)$$

and

$$K^{--}(x,y|x',y') = K(x,y|x',y') - K(x,y|-x',y') \\ - K(x,y|x',-y') + K(x,y|-x',-y') \quad (8d)$$

are defined as well. Equations (7) are enforced for $z = 0$, $x \in (0, L/2)$ and $y \in (0, w/2)$.

Table 1 summarizes the four symmetry cases which are derived as in the foregoing discussion. By means of this table, four integral equation pairs can be constructed in the spirit of (7) by replacing the kernels in (7) with the appropriate kernels from the table and retaining only the non-vanishing terms in the series expansion.

Figure 2 depicts qualitatively the respective modal current distributions for the lowest frequency natural resonance exhibiting each symmetry.

Table 1

COMPATIBLE CURRENT SYMMETRY FEATURES

J _x				J _y				
Sym. w.r.t. x axis	Sym. w.r.t. y axis	Kernel	Compatible Trig. Fns.	Coefs. ≠ 0	Compatible Trig. Fns.	Kernel	Sym. w.r.t. x axis	Sym. w.r.t. y axis
sym	sym	K^{++}	$\cos 2n\phi$	d_{2n+1}^+	$\sin 2n\phi$	K^{--}	anti	anti
sym	anti	K^{+-}	$\cos (2n + 1)\phi$	d_{2n}^+	$\sin (2n + 1)\phi$	K^{-+}	anti	sym
anti	sym	K^{-+}	$\sin (2n + 1)\phi$	d_{2n}^-	$\cos (2n + 1)\phi$	K^{+-}	sym	anti
anti	anti	K^{--}	$\sin 2n\phi$	d_{2n+1}^-	$\cos 2n\phi$	K^{++}	sym	sym

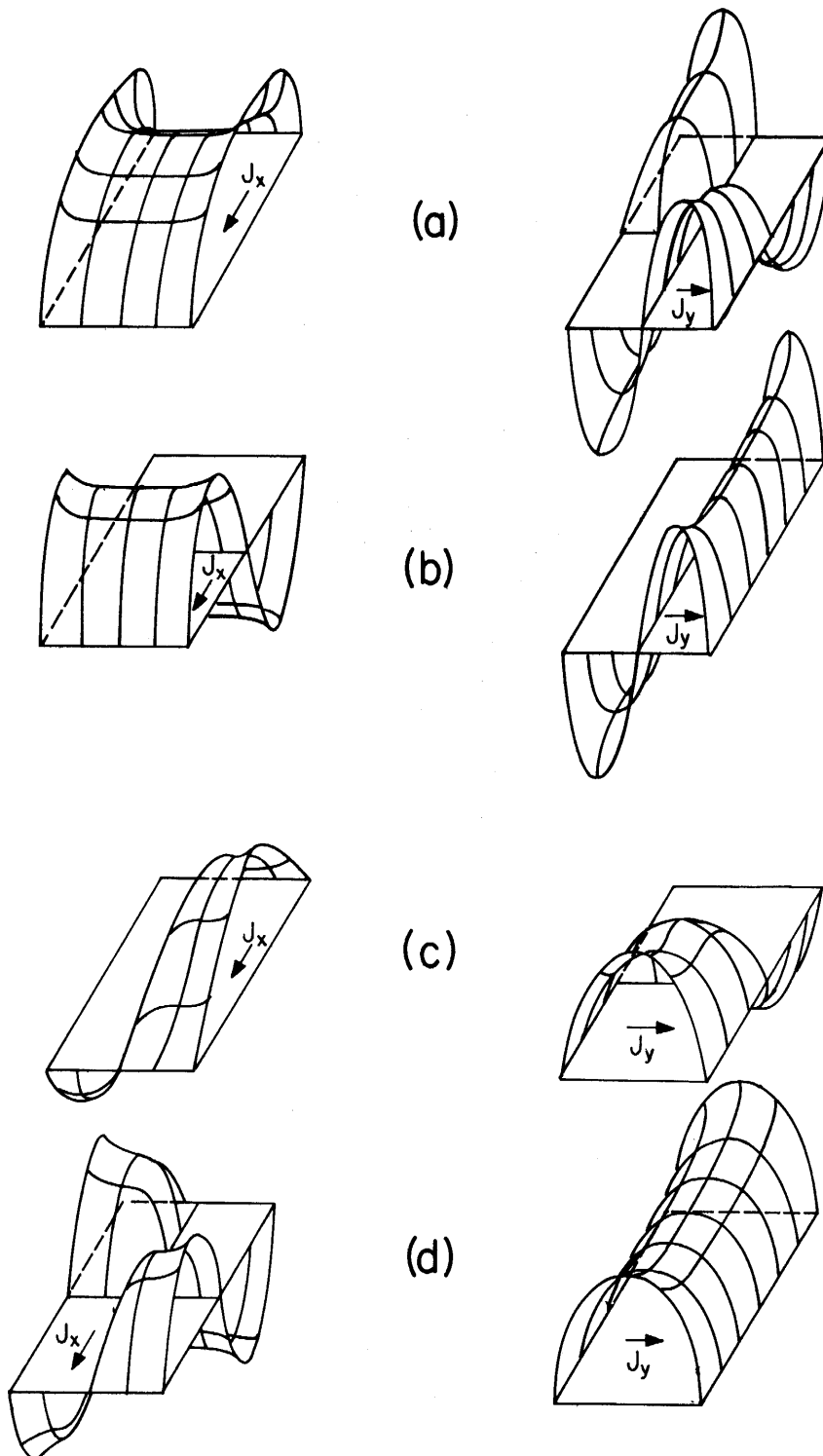


Figure 2. Lowest Order Natural Mode Current Pairs for Each of the Symmetry Cases, a) J_x Symmetric w.r.t. x-Axis and Symmetric w.r.t. y-Axis, b) Symmetric-Antisymmetric, c) Antisymmetric-Symmetric, and d) Antisymmetric-Antisymmetric

SECTION IV
THE NUMERICAL MODEL

The integral equation pair of the form (7) for each of the four symmetry cases can be discretized by the method of moments. In the work reported here, two-dimensional, subsectionally constant expansion functions were used with collocation testing. The zoning scheme is represented in Figure 3.

The unknown currents J_x and J_y were expanded in piecewise constant functions as in (ref. 3) with subsectioning of the form given in Figure 3. Notice that half-width patches are used at the edges of the plate so that match points lie precisely on the edge of the plate. The half-width pulse has proved useful in realizing the actual electrical size of a body in one-dimensional problems (ref. 6). Some numerical experimentation was also done with full-sized pulses on the edges and comparative results are reported in a later section. Some difficulties occur in definition of the edge of the plate in the present formulation because of the presence of two current components which have the asymptotic behavior given in (3). This difficulty is discussed in a later section.

The boundary condition $J_{\text{norm}} = 0$ must be enforced on selected patches at the edge of the plate as discussed in (ref. 3). Concomitantly, only as many d_n^+ 's are retained in the right-hand side summation in (7) as there are current values preassigned to zero. The shaded patches in Figure 3 indicate the selection of patches where a current component is preassigned a zero value. At the corner patch, both components are preassigned zero values.

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6. Butler, C. M., "Integral Equation Solution Methods," in "Wire Antennas and Scatterers," Short Course Notes, University of Mississippi, April 1972. (See also IEEE Trans., v. AP-20, pp. 731-736, 1972.)

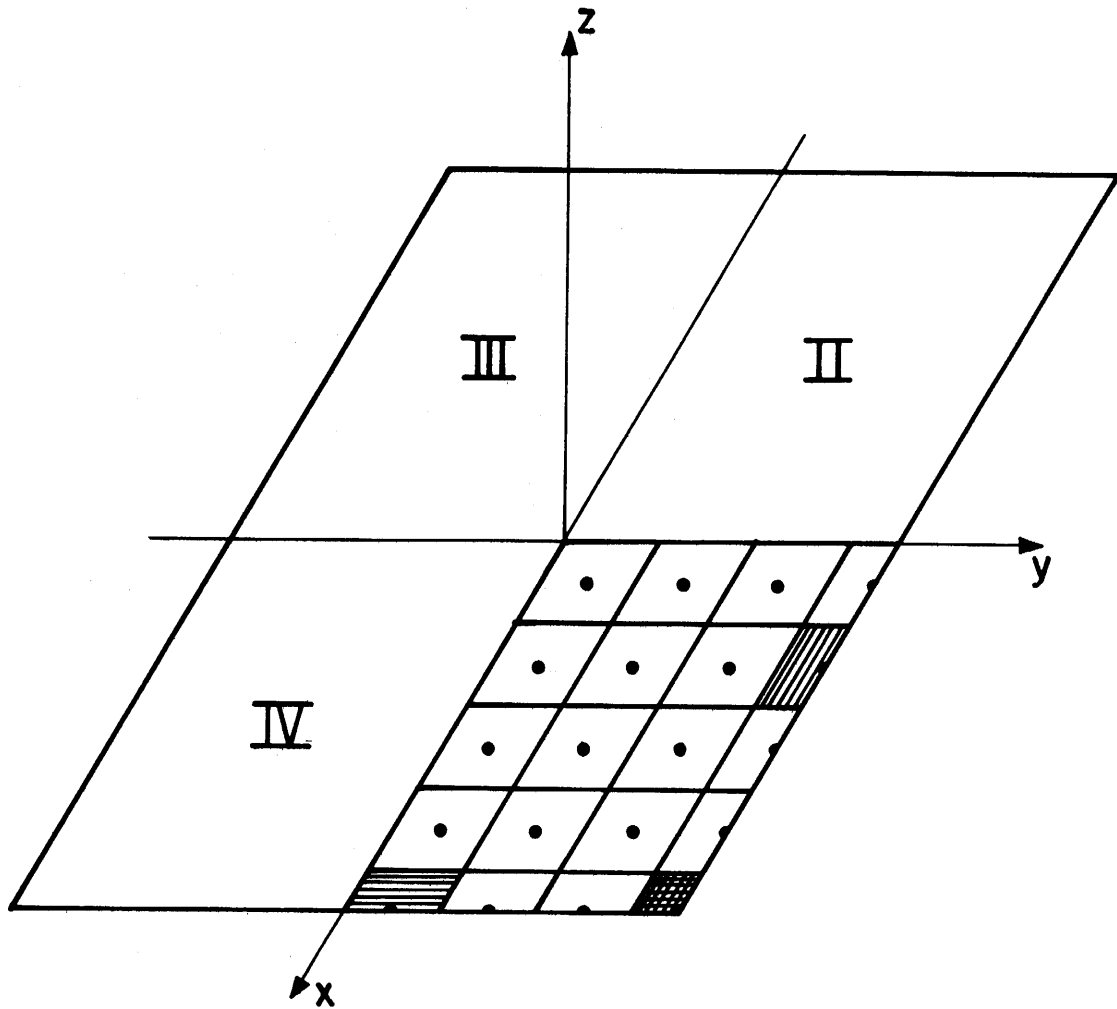


Figure 3. Subsectioning for the Discretization of the Integral Equations

By assigning one match point per expansion patch and by retaining one series expansion term for each current value preassigned in each of the two integral equations, a consistent (i.e. square) system of linear equations results. The truncated summation is taken to the left-hand side so that a homogeneous system results. The matrix organization used to represent these equations is given in Figure 4. Table 2 defines the computer variables noted in Figure 4, primarily for reference purposes in the next section.

The matrix that results is a function of the complex frequency s . A natural resonance occurs when s has a value such that the matrix has a zero determinant; hence, the homogeneous system of equations has a non-trivial solution. The next section explores some algorithmic considerations in the efficient generation and manipulation of the matrix.

$$\begin{bmatrix} \begin{bmatrix} M_x \\ (NI1 \times NJ1) \end{bmatrix} \\ \begin{bmatrix} 0 \\ (NI2 \times NJ1) \end{bmatrix} \end{bmatrix} \begin{bmatrix} 0 \\ (NI1 \times NJ2) \end{bmatrix} \begin{bmatrix} M_\Sigma \\ (NI1 + NI2 \\ \times NPRES) \end{bmatrix} \begin{bmatrix} J_x \\ (NJ1) \\ J_y \\ (NJ2) \\ d \\ (NPRES) \end{bmatrix} = \begin{bmatrix} 0 \end{bmatrix}$$

Figure 4. Organization of the System of Linear Equations

Table 2

MATRIX FORMULATION PARAMETERS

NI1	Number of match points on the zoned quadrant of the plate.
NI2 = NI1	
NPREJ	Number of patches along the $ x = L/2$ edge where J_x is preassigned to zero.
NPREI	Number of patches along the $ y = w/2$ edge where J_y is preassigned to zero.
NJ1 = NI1-NPREJ	Number of unknown current values in J_x expansion.
NJ2 = NI2-NPREI	Number of unknown current values in J_y expansion.
$\left. \begin{array}{l} \text{NJ3} \\ \text{NPRE} \end{array} \right\} = \text{NPREI-NPREJ}$	Number of preassigned current values (Also the number of coefficients retained in summation).

SECTION V

ALGORITHMIC CONSIDERATIONS IN EVALUATING THE SYSTEM DETERMINANT

Some considerations taken into account in generating the system matrix and evaluating its determinant efficiently are discussed in this section. Since these two operations must be repeatedly carried out for many values of s in the course of determining the natural frequencies of the plate, it is essential that clean, efficient computer programming and coding be used so that execution of the program will be affordable. The volume of code in the algorithms is consistently compromised toward a larger size in order to meet the following two time-efficient objectives:

1. Avoidance of calculating the same quantity twice; and
2. Avoidance of logical decisions, particularly those which might be imbedded in loops.

The program is discussed in the context of the following major segments:

1. Computation of an "interaction matrix";
2. Construction of the non-zero submatrices of the system matrix from the interaction matrix;
3. Computation of the series terms' submatrix; and
4. Determinant evaluation.

The major contribution to the elimination of redundant calculations is the one-time computation of an "interaction matrix" which is made up of the individual kernel integral terms from (2) for all argument combinations which occur in the computation. The subsequent program step then picks, by subscript, entries from this matrix and constructs the appropriate kernel from one of equations (8) according to the symmetry conditions being solved. This procedure can be viewed in terms of the layout given in Figure 5a. The terms in the interaction matrix are those evaluated for the match-point as

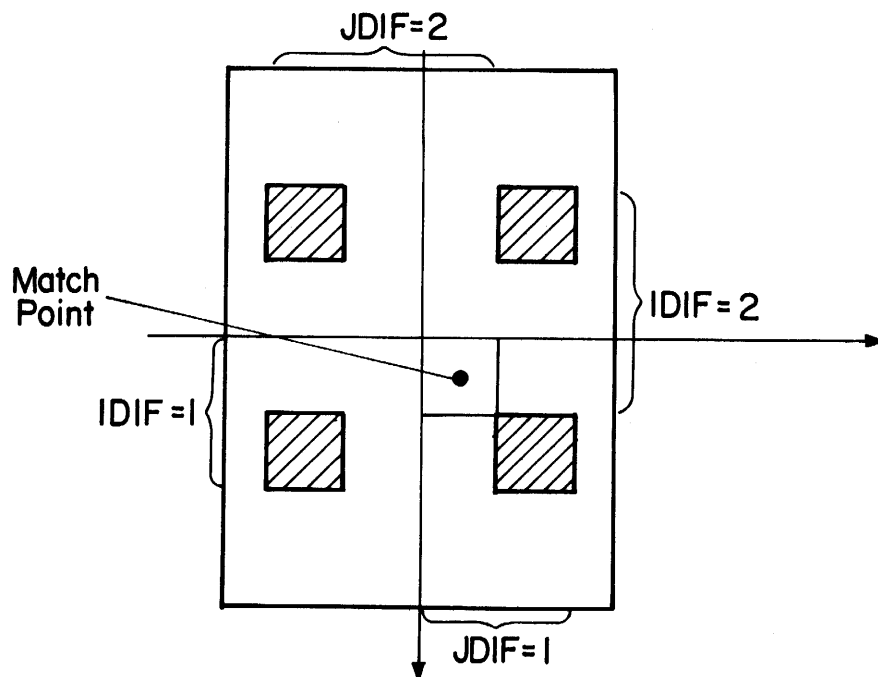
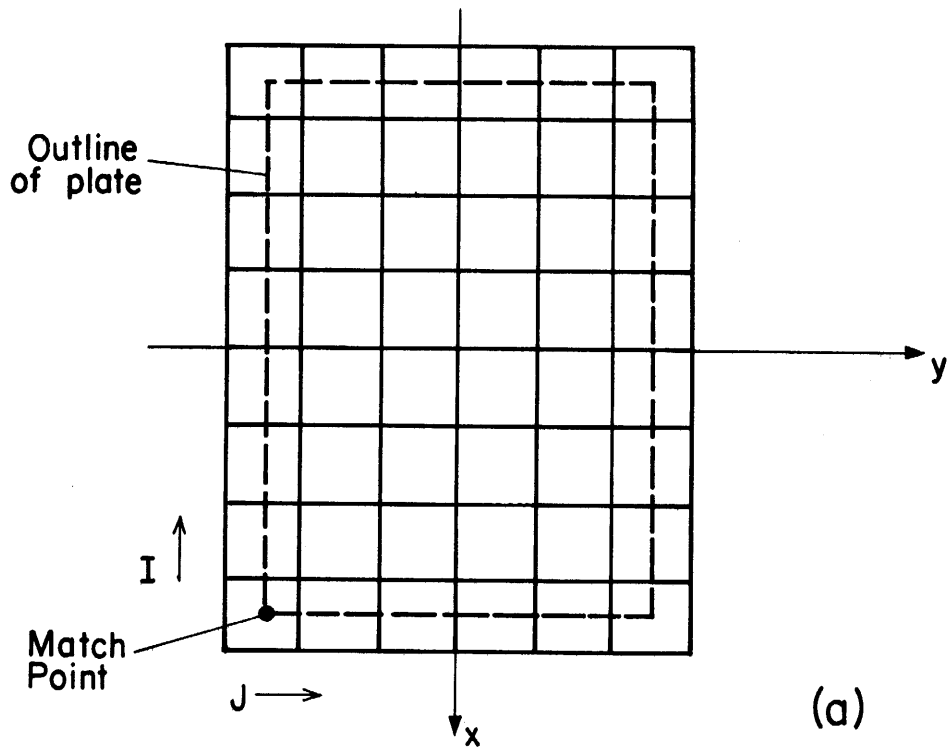


Figure 5. a) Conceptual Zoning for Calculation of the Interaction Matrix (b) Example of the Four Interaction Contributions to a Single Source Term

shown in the lower left with the source patches indexed over the entire plate to generate the matrix. Thus, all geometric relationships which occur in the kernel terms are encompassed in the calculation. Note that all source patches are full patches for this calculation. The effect of half patches at the edges is accounted for by weighting by a factor of 1/2 the edge contributions. The kernel integral appropriate to the symmetry is constructed by summing with correct signs the appropriate elements from the matrix. Figure 5b gives an example of the four source patches entering into one kernel integral.

Differing degrees of sophistication are required in the calculation of the interaction terms depending on the spacing of the patches for which an interaction is being calculated. For the self patch, i.e., the patch in which the match-point resides, the integration of the kernel must be performed analytically because of the integrable singularity in the kernel there. Appendix A gives a series approximation to this integral. The result in Appendix A is evaluated directly in the program. For the patches adjacent to the patch containing the match point, the kernel is a rapidly varying but well-behaved function. The integration over these patches is evaluated numerically by a polynomial approximation. For patches further separated, the kernel is slowly varying and the integral is evaluated approximately as the product of the value of the kernel at the center of the patch and the area of the patch.

Some minor time economy is achieved in the matrix filling algorithm, which is a four-dimensional loop. The economy is found in the form of decision-free indexing, that is, the source contributions from interior patches, from $|x| = L/2$ edge patches, from $|y| = w/2$ edge patches, and from corners take on different forms. Rather than index over all source patches

with logical decisions implemented to discriminate among the four cases above, four different loops are used.

The computation of the series term submatrix is relatively straightforward. Because the Bessel-trigonometric products appear in two terms each, they are all precalculated and stored in a vector. The individual terms are then constructed from them.

The determinant evaluation profits significantly from an exploitation of the sparceness of the matrix. Either of two approaches may be taken to the sparse matrix manipulations. One is to separate the matrix algebraically and calculate an inverse as a composite of inverses of terms involving the submatrices. The alternative approach is to attack the matrix directly with Gaussian elimination, and to exploit the sparceness directly in the algorithm. The latter approach was chosen for the present purpose because it is judged to be slightly faster computationally and because in order to determine natural mode solutions for the SEM formulation, the homogeneous system of equations occurring at a pole must be backsolved. The algorithm resulting from the second approach is described in Appendix B.

The determinant evaluation routine itself appears in Appendix C as the function routine CPLATE.

SECTION VI

NUMERICAL CHECKS ON THE ACCURACY OF THE POLES

The results of some numerical checks on the accuracy of the pole location as determined from the numerical model described in Sections II through V are reported. It is shown that the model predicts well the poles for narrow strips possessing essentially thin scatterer resonance properties. Difficulties occur, however, in obtaining self-consistent results under different zone sizes for plates with larger aspect ratios. It is conjectured that the difficulty occurs because the subsectionally constant current representation is inadequate to represent the correct edge behavior for the currents—particularly the singular behavior for the parallel component. The rationale behind this conjecture is discussed.

Initial tests on the accuracy of the model were made for a rectangular strip with a shape ratio $w/L = 1/10$. Such a strip has an approximate equivalent dipole whose diameter-to-length ratio is $1/10\pi$.

Figure 6 gives the results of pole determinations for the first two poles for various numbers of pulses in the expansion of the current and for two different treatments of the edge pulse. The strip was zoned with one pulse across a quadrant. The numbers indicated in the figure are the numbers of pulses along the longitudinal direction of a quadrant. The "half-pulse" results are those obtained by the zone scheme described in Section IV where half-width pulses are placed at the edge so that match points fall at the edge. The "full-pulse" results are those obtained by zoning the plate with equal-sized pulses over the entire quadrant. In the latter case an a posteriori adjustment is made in the data correcting the length of the strip such that the end of the corrected strip lies at the end match-point.

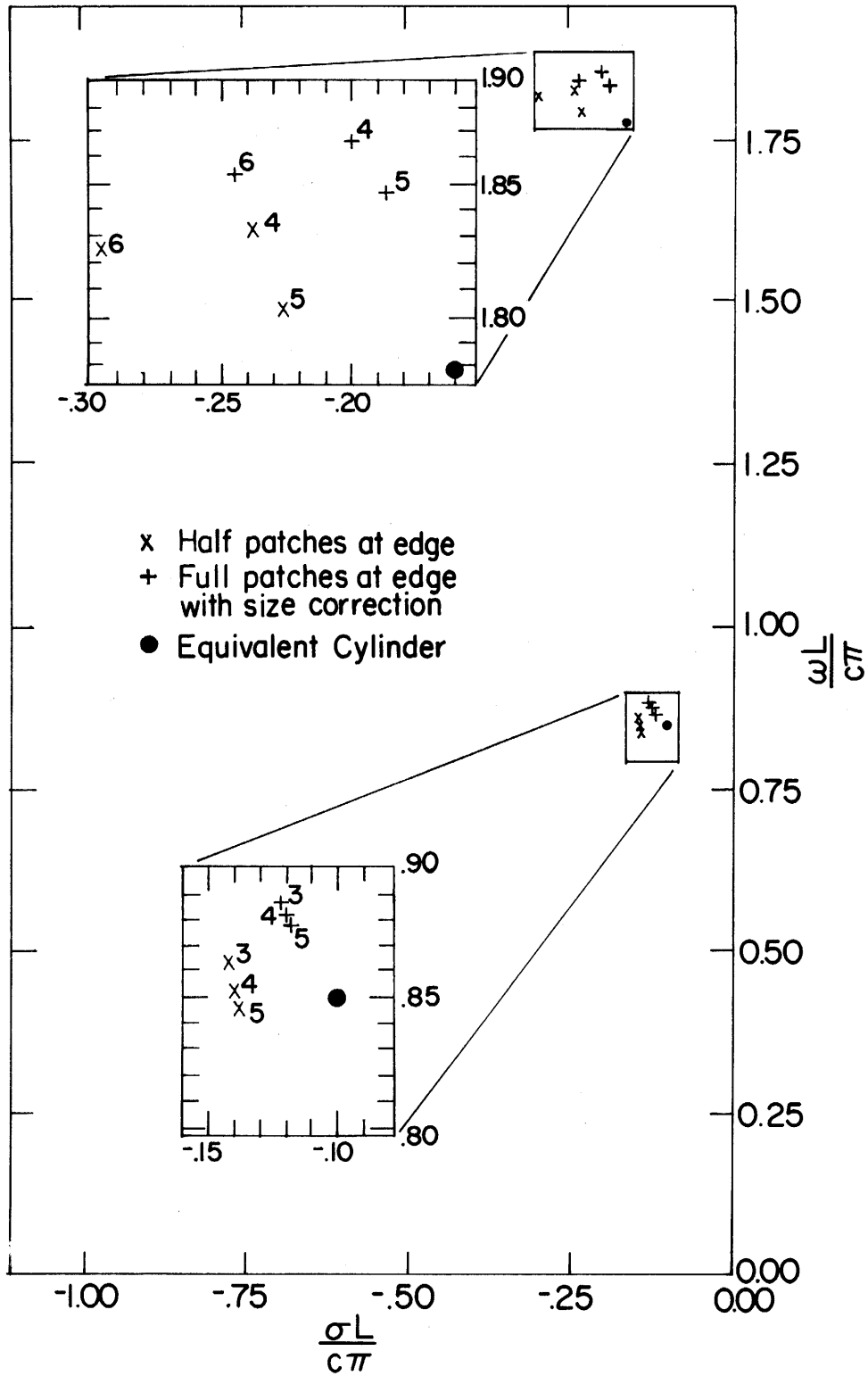


Figure 6. Calculated Pole Locations for Thin-Strip for Varying Numbers of Zones in the x-Direction and Different Edge Treatments (Cylinder Results from Ref. 6)

It is seen that the differences are small both for varying order and increasing pulse density. The $NX = 6$ results for the second pole show some departure from the trend established by the results for $NX = 4$ and $NX = 5$. This is attributable to the fact that the matrix is on the brink of numerical instability for $NX = 6$. The results for $NX = 7$, which are not shown, are observed to be meaningless because of the instability manifested.

For comparison purposes, the first two poles for an equivalent cylinder (one whose circumference equals the strip width) are given as found in ref. 7. These results are judged reliable inasmuch as they have been cross-checked among several integral equation formulations and for several method-of-moments schemes. The equivalent radius taken is, of course, an approximation. It is seen that the half-pulse solutions compare slightly more favorably with the cylinder results. Because of this, and moreover, because the a posteriori data adjustment is avoided with the half-pulse scheme, it was used consistently in the remaining solutions.

The stable convergence properties of the numerical model exhibited for the thin-strip solution are not manifested for higher aspect ratios. The reason for the difference is that the strip is essentially a one-dimensional problem and the transverse (y-directed) component of current has little effect on the dominant longitudinal current. For wider structures the coupling is significant and inadequacies in the modeling of the singularities of the currents produce inaccuracies which are highly sensitive to zoning.

Figure 7 shows the results obtained for a pole trajectory as a function of the shape factor w/L where the zoning in the y-direction was adjusted

-
7. Umashankar, K. R., "Transient Scattering by a Thin Wire in Free Space and Above Ground Plane Using the Singularity Expansion Method," Interaction Note 236, August 1974. (See also F. M. Tesche, Interaction Note 102, April 1972.)

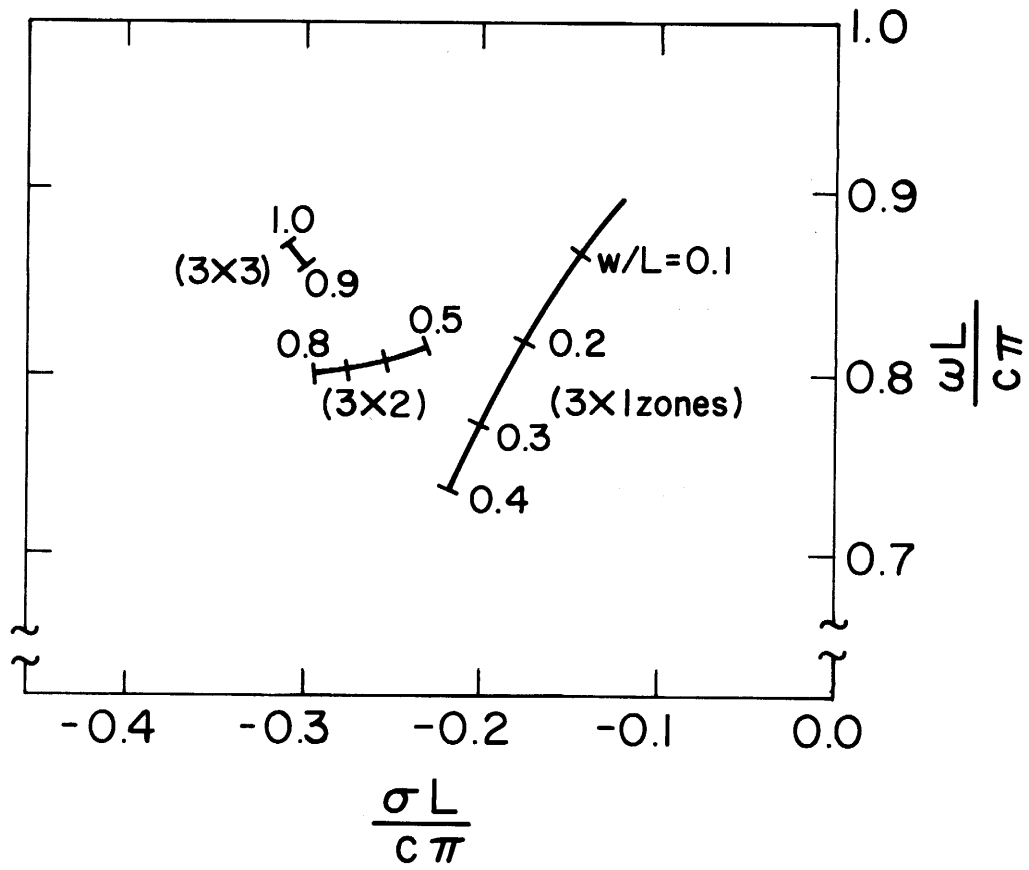


Figure 7. Computer Pole Trajectory Under Varying w/L with Zoning Changes

according to the value of w . It is evident that the solutions are unstable with respect to the zoning on the plate. Attempts to increase the number of zones significantly to improve upon the situation resulted in numerical instabilities in the matrix which cause the pole search iteration to fail.

The reason for the difficulty manifested in Figure 6 is believed to lie in the way that the edge of the plate is defined with the piecewise constant current expansion. Consider the characteristics of the two current components along a line traversing the plate in the y -direction as shown in Figure 8. The correct edge behavior at $|y| = w/2$ is that given in equations (3). The zoning scheme, however, forces $J_x(x, \pm w/2)$ to take a finite value. The current extrapolates to a singular point for some $y > w/2$, i.e., the numerical model represents a plate whose width is greater than w .

If the number of transverse zones is increased as indicated by the dashed curve in Figure 8, the steepness of the edge behavior of J_x is increased, and the extrapolation is characteristic of a narrower plate as compared to the first case. This narrowing of the effective width in the model is reflected in an increased Q (resonance quality factor) as the jumps in Figure 7 indicate.

One is tempted to conclude that a full-width pulse at the edge is a cure for the problem since the point at which the pulse current is defined is shifted relative to the edge as zoning is changed with full-width pulses. The effect of this procedure is to transfer the problem from component of current whose behavior is singular at the edge to the component which goes to zero. With full pulses at the edges, the normal component of current would go to zero interior to the plate rather than at the edge as it properly should.

An approach which is potentially a remedy for this difficulty is discussed in the conclusions.

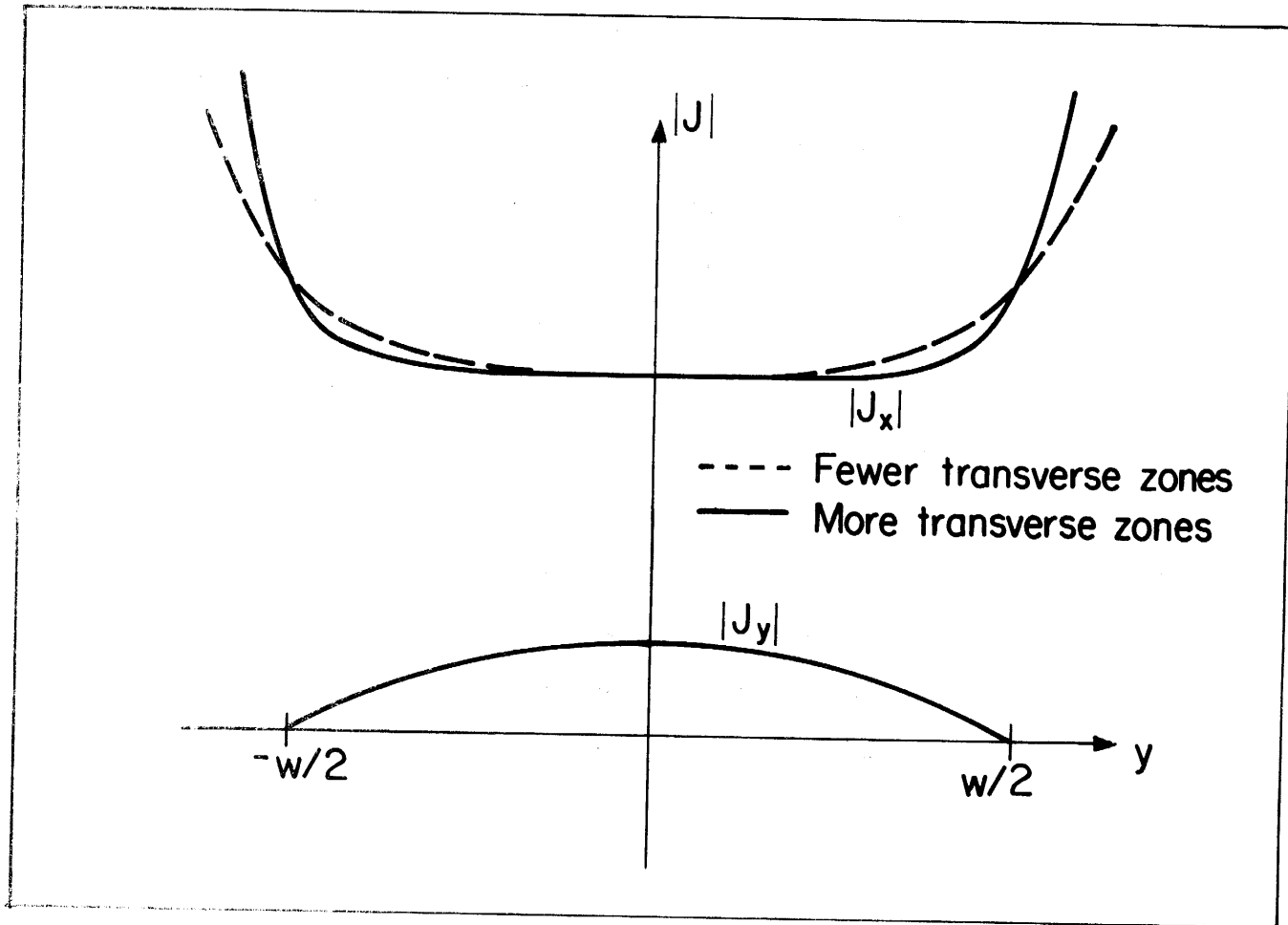


Figure 8. Behavior of Singular Component of Current at the Edge Under Change in Transverse Zoning

SECTION VII

POLE TRAJECTORIES AS A FUNCTION OF SHAPE RATIO

Figure 9 gives the results obtained for pole location for the lowest order pole of each of the symmetries as a function of w/L . Clearly, as the previous section indicates, the results cannot be taken as the correct results for the plate. However, the zoning was fixed for all w/L and the results are expected to reflect the proper trends for these trajectories.

The ++ and +- modes are in essence dipole modes, and their poles show the general lowering of Q as w/L increases. (The ++ indicates that the J_x component is symmetric both with respect to the x and y axes - see Table I.) The -- mode can be thought of as a dipole mode in the transverse direction. As a result it shows a strong frequency dependence on the transverse dimension w . When $w/L = 1$, the -- mode is identical to the ++ mode rotated spatially 90 degrees. Consequently, the two trajectories coalesce as $w/L \rightarrow 1$, when the zoning is 5×5 so as to preserve symmetry in the numerical mode. The failure of the 5×3 zone case is due to the reasons outlined in the previous section.

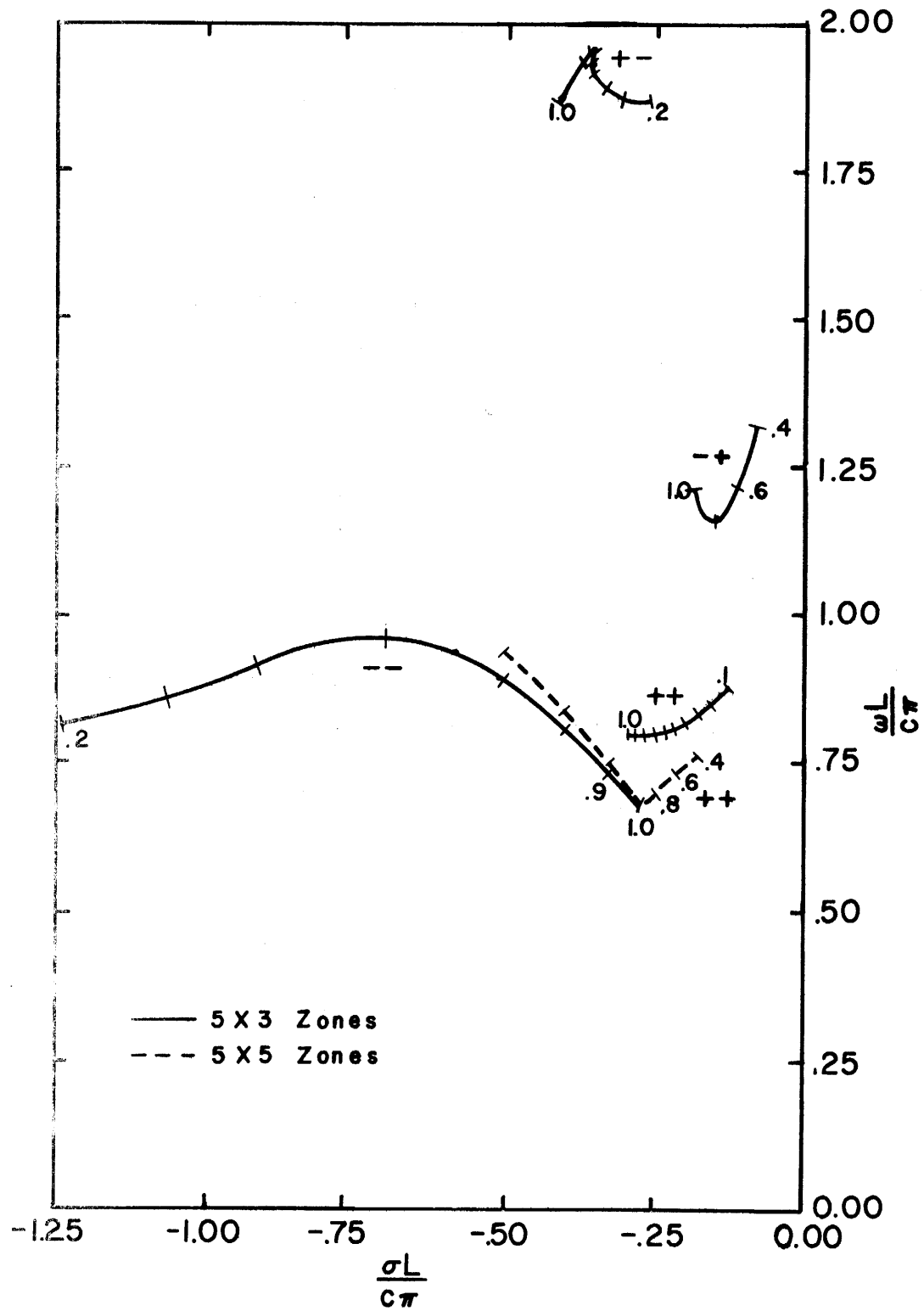


Figure 9. Pole Trajectories as Computed with Zoning Fixed

SECTION VIII

CONCLUSIONS

The application of SEM to the equivalent problems of the perfectly conducting rectangular plate and the rectangular aperture in a perfectly conducting screen has been conducted with partial success. The applicability of SEM and the computational feasibility of determining SEM quantities are demonstrated. It is further demonstrated that by careful program construction, the computational costs of numerical treatment of two-dimensional problems can be made quite reasonable. The cost of generating a matrix and calculating its determinant by the methods discussed herein is less than 10 cents for each value of s .

Difficulties arise in the subsectionally constant current zoning because of a failure to properly model the edge conditions. Whereas Rahmat-Samii and Mittra (ref. 3) obtained good radar cross-section results with such a zoning scheme, the pole locations are highly sensitive to the zoning. Radar cross-section is a far-field quantity, and the integrated effects of the errors are small there. The pole locations, on the other hand, depend strongly on the dimensions of the structure, and it is evident that the plate size must be brought to bear in a precise fashion for the pole locations to be correct.

The edge problem can be remedied by imposing the edge conditions (3) directly in the basis set elements for edge zones. Davis has done this successfully for the circumferential component of current on a thick cylindrical scatterer (ref. 8). The circumferential current

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8. Davis, W. A., "Numerical Solutions to the Problem of Electromagnetic Radiation and Scattering by a Finite Cylinder," Ph.D. Thesis, University of Illinois, 1974.

component is singular at the ends of the cylinder. The introduction of the singular basis element will produce a significant complication to the problem in that a second singularity, that of the current, must be integrated analytically in order to implement the model with edge conditions imposed. An independent check on program accuracy is dictated for a problem of this scope before proceeding with the edge condition approach.

APPENDIX A

THE SELF-PATCH INTEGRATION

The term for the interaction matrix for IDIF = JDIF = 0, i.e., where the match point lies at the center of the source patch, can be written

$$I_s = 4 \int_0^{\Delta x/2} \int_0^{\Delta y/2} K(0,0|x',y') dx' dy' \quad (A1)$$

This presumes a unit amplitude expansion pulse over the patch whose dimensions are Δx and Δy . The symmetry of the kernel with respect to both x' and y' is employed in the forming of (A1). This integral can be transformed to polar coordinates as

$$\begin{aligned} I_s &= 4 \left\{ \int_{\phi=0}^{\tan^{-1} \frac{\Delta y}{\Delta x}} \int_{\rho=0}^{\frac{\Delta x}{2 \cos \phi}} \exp[-s\rho/c] d\rho d\phi \right. \\ &\quad \left. + \int_{\phi=\tan^{-1} \frac{\Delta y}{\Delta x}}^{\pi/2} \int_{\rho=0}^{\frac{\Delta x}{2 \cos \phi}} \exp[-s\rho/c] d\rho d\phi \right\} \\ &= -\frac{4c}{s} \left\{ \int_{\phi=0}^{\tan^{-1} \frac{\Delta y}{\Delta x}} [\exp(-s\Delta x \sec \phi/2c) - 1] d\phi \right. \\ &\quad \left. + \int_{\phi=\tan^{-1} \frac{\Delta y}{\Delta x}}^{\pi/2} [\exp(-s\Delta y \csc \phi/2c) - 1] d\phi \right\} \quad (A2) \end{aligned}$$

If the exponential functions in the integrand are then expanded in McLaurin series, the resulting terms of powers of secants and cosecants possess tabulated integrals. The result for three terms retained in the series is

$$\begin{aligned}
I_s \approx & -\frac{4c}{s} \left\{ -\frac{s\Delta x}{2c} \cdot \frac{1}{2} \ln q_y + \frac{1}{2} \left(\frac{s\Delta x}{2c}\right)^2 \frac{\Delta y}{\Delta x} \right. \\
& - \frac{1}{6} \left(\frac{s\Delta x}{2c}\right)^3 \frac{\Delta x(\Delta x^2 + \Delta y^2)^{1/2}}{2\Delta y^2} - \frac{s\Delta y}{2c} \frac{1}{2} \ln q_x \\
& \left. + \frac{1}{2} \left(\frac{s\Delta y}{2c}\right)^2 \frac{\Delta x}{\Delta y} - \frac{1}{6} \left(\frac{s\Delta y}{2c}\right)^3 \frac{\Delta y(\Delta x^2 + \Delta y^2)^{1/2}}{2\Delta y^2} \right\} \quad (A3)
\end{aligned}$$

where

$$q \begin{matrix} (x) \\ (y) \end{matrix} = \frac{[(\Delta x^2 + \Delta y^2)^{1/2} + (\Delta x)]}{[(\Delta x^2 + \Delta y^2)^{1/2} - (\Delta x)]}$$

APPENDIX B

THE SPARSE MATRIX ALGORITHMS

1. Introduction

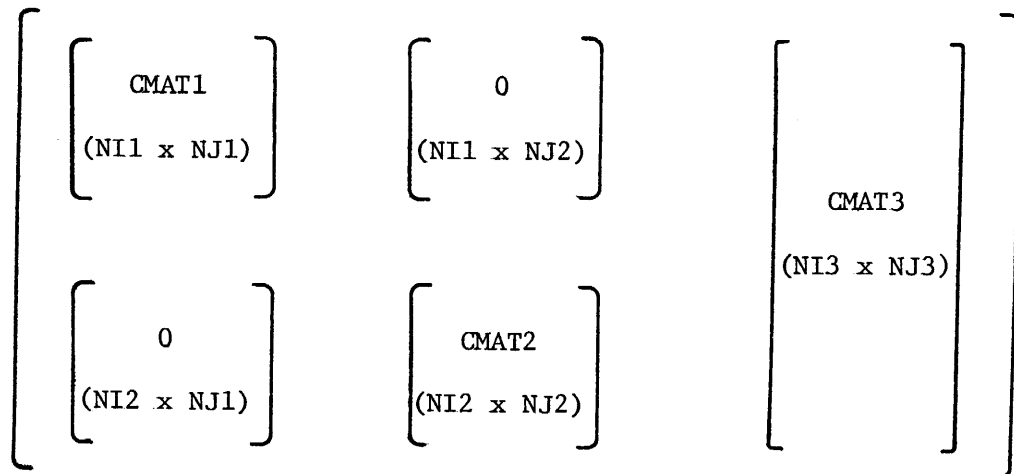
This Appendix describes the algorithmic approach to minimize the computation time involved in Gaussian elimination triangularization of systems of matrix equations which are "sparsely coupled." The term "sparsely coupled" as applied in this Appendix implies the matrix equation form given in (B1).

$$[M] [X] = \begin{bmatrix} M_1 & 0 & \\ & & M_3 \\ 0 & M_2 & \end{bmatrix} \begin{bmatrix} X \\ \\ \end{bmatrix} = \begin{bmatrix} B \\ \\ \end{bmatrix} \quad (B1)$$

It is clear that in this form the sole coupling between the "upper" and "lower" systems of equations is contained in the matrix M_2 . Generally, the number of columns in M_2 is small compared with the order of the overall system.

An algebraic approach to exploiting the sparceness of (B1) results in solutions which are given in terms of the several submatrices and their inverses. (See, for example, ref. 9.) It is well-known, however, that it is sufficient for the purposes of determinant calculation and equation solution to triangularize the composite matrix in (B1). The triangularization process involves fewer operations than the diagonalization necessary for the calculation of an inverse.

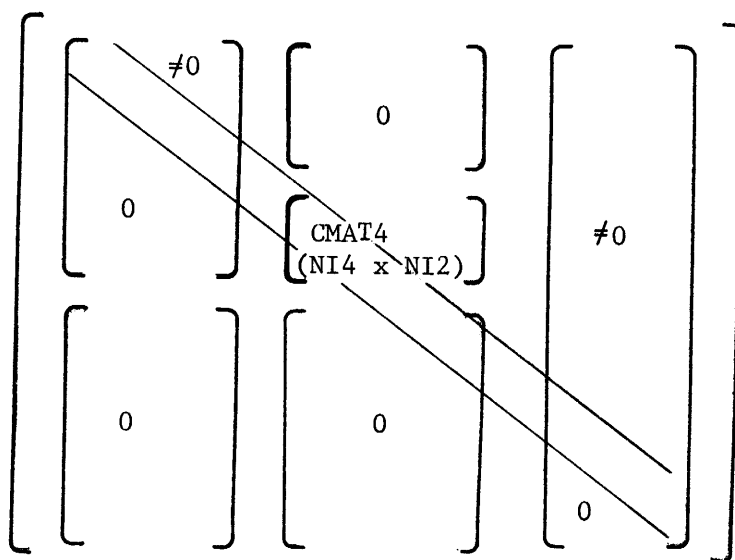
-
9. Dunaway, O. C., "A Numerical Solution for the Distribution of Time-Harmonic Electromagnetic Fields in an Arbitrary Shaped Aperture in a Ground Screen," M.S. Thesis, University of Mississippi, 1974.



$$\text{NI3} = \text{NI1} + \text{NI2}$$

$$\text{NJ3} = \text{NI3} - \text{NJ1} - \text{NJ2}$$

(a)



$$\text{NI4} = \text{MAX} (\text{NI1} - \text{NJ1}, 0)$$

(b)

Figure B1. Submatrix Organization for the Sparse Matrix Algorithms, a) the Original Matrix, and b) Triangularized Form with the Generated CMAT4

This Appendix describes an algorithmic exploitation of the sparseness of the composite matrix in (B1). That is, a numerical process is given whereby only the non-zero subelements are stored and operated on, with the computational operations being those which render the composite matrix M upper triangular. The determinant of the composite matrix results directly from this triangularization. A solution for X in (B1) requires a backsolving process involving the triangularized form of M and a vector resulting from applying the elimination operations to the vector B . Routines to perform these operations are given as well.

Appendix C gives listings of the routines built on this algorithm. The triangularization routine is named SPRHOM. The backsolving procedure is performed by the entry HOMSLV to the routine SPRSLV. (The name entry SPRSLV backsolves an inhomogeneous system and is not used for present purposes.)

2. The Algorithm

The routine SPRHOM is simply an implementation of the Gaussian elimination procedure with maximum pivot selection applied to the composite matrix M in (B1) viewed as a whole. The sparseness of M is exploited by storing only the non-zero submatrices in (B1) and skipping the operations involving zero elements. The result is a substantial saving in both storage and computation time.

The forms of the input and of the end product for the triangularization are given in Figure (B1) with the sizes of the respective submatrices defined. It is recalled that the fundamental process in the Gaussian elimination procedure is the elimination of all or part of the elements of a column of a matrix with respect to a pivot element, commonly the element of greatest magnitude in the column. That is, for a column under process, the row

containing the main diagonal element of the matrix which falls in that column. All or part of the elements not on the main diagonal are "eliminated" or made zero by subtraction of some multiple of the row containing the column maximum. In the triangularization procedure, the part of the column comprising elements lying below the main diagonal after row exchange are eliminated. If the matrix is a part of a system of equations with non-zero right-hand side, the row operations of exchange and subtraction of a constant multiple of another row must be performed on the corresponding elements of the right-hand side vector as well.

The algorithm of the routine SPRHOM operates according to the Gaussian elimination procedure described above. However, the three submatrices CMAT1, CMAT2, and CMAT3 are stored individually. In addition, the routine generates a submatrix CMAT4 in the course of selecting pivots for the columns contained in CMAT2. Further, the "elimination" of elements of submatrices that are zero a priori is skipped. The result is significant storage and time economy.

In order to minimize logic decisions in the routine, it is organized to operate sequentially through the partitioned matrix. The steps are as follows (it is convenient to follow the thinking of these steps by tracing the location diagonal of the composite with the aid of Table B1):

- a. Perform conventional Gaussian elimination to zero the elements $CMAT1(I,J)$ for $I > J$, i.e., the elements below the main diagonal of M. Choose maximum pivot elements in conventional fashion. Carry row operations into CMAT3.
- b. Create CMAT4 by row swapping with CMAT2 so as to choose maximum pivot elements. Perform elimination to zero CMAT4 elements for $I > J$ and the entire column of CMAT2. The number of rows created in CMAT2 is $NI4 = NI1 - NJ1$, the amount by which CMAT1 is over-square. Carry row operations across into CMAT3.

- c. Choose maximum pivot rows in columns of CMAT2 with $J > NI4$ and swap with rows given by $I = J - NI4$ (the rows containing the Jth column diagonal element of the composite). Conduct elimination to zero elements with $I > J + NI4$. Carry row operations into CMAT3.
- d. Conduct conventional pivot selection and elimination on CMAT3 to zero elements CMAT3(I, J) with $I > J + NJ1 + NJ2$.

In each column elimination operation, the pivot value is multiplied into a product accumulator to produce a value for the determinant of the composite matrix. The sign of this product is changed at each row swap in accord with the rates of matrix algebra row operations.

The backsolving routine SPRSLV with its entry HOMSLV operate in a straightforward manner on the submatrices as reduced by SPRHOM. Details are omitted here, but the routines may be easily followed in a row-by-row flow from the bottom of the composite matrix, if one keeps in mind the row index relations of column 4 of Table B1. The entry HOMSLV assumes a zero determinant value resulted (approximately) from SPRHOM and the last element of the solution vector is picked as unity. (The zero determinant results from a zero falling at the last diagonal location in maximum pivoting Gaussian elimination.) The remainder of the homogeneous solution vector is backsolved conventionally relative to this last element. The vector is then renormalized so that its maximum element is unity.

Table B1

PRIMARY INDEXING QUANTITIES IN THE ALGORITHM

Submatrix	Size of ¹ Submatrix	Indices of Main ¹ Diag. of Compos.	Relative Row ² Index of CMAT3 and CRHS
CMAT1	NI1 x NI2	(J,J)	I3 = I
CMAT4	NI1 - NJ1 x NJ2 (can be null)	(J,J)	I3 = I + NJ1
CMAT2	NI2 x NJ2	(J - (NI1 - NJ1), J)	I3 = I + NI1
CMAT3	NI1 + NI2 x NI1 + NI2 - NJ1 - NJ2	(J + NJ1 + NJ2, J)	I3 = I3

1. Quantities given in terms of input parms. to the routine. Related internal quantities are given in Figure B1.
2. Relative to I, the row index of the submatrix in question.

APPENDIX C
PROGRAM LISTINGS

All code compileable on IBM OS/360 and OS/370 FORTRAN levels G or H. The routine ZANLYT and its service routine UERTST is taken from the program library FORTUOI made available by the Computer Services Office, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801. The routines BSLJZ and BSCJZ are taken from the International Mathematical and Statistical Library (IMSL). They may not be reproduced apart from this application program package. The IMSL library is available by subscription from IMSL, Inc., 6100 Hillcroft, Suite 510, Houston, Texas 77036.

```

C      POLE SEARCH PROGRAM FOR S E M FORMULATION OF THIN-PLATE SCATTERER      00010
C      BY L W PEARSON 8/74                                                    00020
C                                                                              00030
      IMPLICIT REAL*8(A,B,D-H,O-Z),COMPLEX*16(C)                             00040
      COMMON /GE3M/ XSYM,YSYM,W,NX,NY,IPREAS(10),JPREAS(10),NPREI,NPREJ      00050
      INTEGER MES(4,2) /'SYMM','ETRI','C',' ','ANTI','SYMM','ETRI','C' /    00060
      DATA C/(3.D08,0.D0) /,PLUS/'+' /,PI/3.141592653589793/              00070
      DATA HX/'X' /,HY/'Y' /                                                00080
      EXTERNAL CPLATE                                                         00090
      DIMENSION CX(20),INFER(20)                                             00100
      LOGICAL LAUTO                                                           00110
100    READ(5,1,END=999) XSYM,YSYM,NX,NY,WO,WS,W,CSUNOR,LAUTO                00120
1      FORMAT(2A1,2X,2I3,5F10.4,T80,L1)                                       00130
      IMX=1                                                                    00140
      IMY=1                                                                    00150
      IF(XSYM.NE.PLUS) IMX=2                                                  00160
      IF(YSYM.NE.PLUS) IMY=2                                                  00170
      NW=(W-WO)/WS                                                            00180
      IF(NW.GT.0) GO TO 105                                                    00190
      NW=-NW                                                                    00200
      WS=-WS                                                                    00210
105    IF(WS*NW.LT.W-WO) NW=NW+1                                             00220
      DO 200 IW=1,NW                                                           00230
      W=WO+(IW-1)*WS                                                         00240
      IF(.NOT.LAUTO) GO TO 140                                                00250
                                                                              00260
                                                                              SKIP PAST AUTO ZONING
                                                                              00270
      ROUTINE TO DETERMINE NO OF EXPANSION PULSES BASED ON ELECTRICAL        00280
      SIZE OF PLATE                                                           00290
                                                                              00300
      TESTWV=.1885D10/DABS(DIMAG(CSUNOR))                                     00310
                                                                              00320
      ////////////////////////////////////                                       00330
                                                                              00340
      NPPWVL=6                                                                00350
                                                                              00360
      ////////////////////////////////////                                       00370
                                                                              00380
      FLENX=1/TESTWV                                                         00390
      NX=IDINT(FLENX*NPPWVL)                                                  00400
      IF(DFLOAT(NX).LT.FLENX*NPPWVL) NX=NX+1                                00410
      FLENY=W/TESTWV                                                         00420
      NY=IDINT(FLENY*NPPWVL)                                                  00430
      IF(DFLOAT(NY).LT.FLENY*NPPWVL) NY=NY+1                                00440
      NX=MINO(NX,5)                                                           00450
      NY=MINO(NY,5)                                                           00460
                                                                              00470
      BEGIN SETUP FOR ALTERNATE EDGE PATCH PREASSIGNMENT                    00480
                                                                              00490
140    NPREI=(NX+2)/3                                                         00500
      NPREJ=(NY+2)/3                                                         00510
      IF(NX-2*NPREI+2.LE.1.AND.NPREI.GT.1) NPREI=NPREI-1                   00520
      IF(NY-2*NPREJ+2.LE.1.AND.NPREJ.GT.1) NPREJ=NPREJ-1                   00530
      DO 110 I=1,NPREI                                                         00540
      IPREAS(NPREI+1-I)=NX-3*I+3                                             00550
110    CONTINUE                                                               00560
      DO 120 J=1,NPREJ                                                         00570
      JPREAS(NPREJ+1-J)=NY-3*J+3                                             00580
120    CONTINUE                                                               00590
                                                                              00600
                                                                              LOCATIONS WHERE X-DIRECTED CURREN
                                                                              T IS SET TO ZERO GIVEN BY SUBSCRI
                                                                              00610

```

C
C
C

PTS (NX,JPRES) AND Y-DIRECTED BY
(IPREAS,NY)

00620
00630
00640
00650
00660
00670
00680
00690
00700
00710
00720
00730
00740
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00790
00800
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00820
00830
00840

WRITE(6,2) W,CSUNDR
2 FORMAT('ENTER ITERATION',/, 'OSHAPE RATIO =',F5.3,5X,
1 'STARTING FREQ =',2D12.4)
WRITE(6,3)
3 FORMAT('0',10X,'CUR SYMMETRY',6X,'PULSES',3X,'PREASSIGN LOC''NS')
WRITE(6,4) HX,(MES(I,IMX),I=1,4),NX,(IPREAS(J),J=1,NPREI)
4 FORMAT(' ',A1,'-DIR',5X,4A4,I6,5X,10I3)
WRITE(6,4) HY,(MES(I,IMY),I=1,4),NY,(JPRES(J),J=1,NPREJ)
WRITE(6,5)
5 FORMAT('0',11X,'COMPLEX FREQ',17X,'DETERMINANT')
CX(1)=CSUNDR
CALL ZANLYT(CPLATE,1.0-50,4,0,1,1,CX,100,INFER,IER)
WRITE(6,6) CX(1)
6 FORMAT('0RETURN FROM ITERATION',/, 'OPOLE LOC''N',2E12.4)
CALL MODE
CSUNDR=CX(1)
200 CONTINUE
GO TO 100
999 STOP
END

	SUBROUTINE MODE	00850
	IMPLICIT REAL*8(A,B,D-H,O-Z),COMPLEX*16(C)	00860
	COMMON /MAT/ CMATX(25,25),CMATY(25,25),CHOM(50,10),CMAT4(10,25),	00870
	INPTCHS,NDIM1,NDIMCI,NDIMCJ,NORD	00880
	COMMON /GEOM/ XSYM,YSYM,W,NX,NY,IPREAS(10),JPREAS(10),NPREI,NPREJ	00890
	DIMENSION CPRX(5,5),CPRY(5,5)	00900
	DIMENSION CJ(50)	00910
	NPRE=NPREI+NPREJ	00920
	NPREIM=NPREI-1	00930
	NPREJM=NPREJ-1	00940
	CALL HOMSLV(CMATX,NPTCHS,NPTCHS-NPREJ,NDIM1,NDIM1,	00950
1	CMATY,NPTCHS,NPTCHS-NPREI,NDIM1,NDIM1,	00960
2	CHOM,NDIMCI,NDIMCJ,CMAT4,NDIMCJ,NDIM1,CJ,NORD)	00970
	NXMI=NX-1	00980
	NYMI=NY-1	00990
	NSUBS=0	01000
	DO 470 JS=1,NY	01010
	DO 450 IS=1,NXMI	01020
	J=(JS-1)*NX+IS	01030
	CPRX(IS,JS)=CJ(J-NSUBS)	01040
	JM=J-NSUBS	01050
450	CONTINUE	01060
	J=JS*NX	01070
	IF(JS.NE.JPREAS(NSUBS+1)) GO TO 460	01080
	NSUBS=MINO(NSUBS+1,NPREJM)	01090
	CPRX(NX,JS)=(0.,0.)	01100
	GO TO 470	01110
460	CPRX(NX,JS)=CJ(J-NSUBS)	01120
470	CONTINUE	01130
	DO 500 IS=1,NX	01140
	DO 500 JS=1,NYMI	01150
	J=(JS-1)*NX+IS	01160
	CPRY(IS,JS)=CJ(NPTCHS-NPREJ+J)	01170
500	CONTINUE	01180
	NSUBS=0	01190
	DO 530 IS=1,NX	01200
	J=NYMI*NX+IS	01210
	IF(IS.NE.IPREAS(NSUBS+1)) GO TO 510	01220
	CPRY(IS,NY)=(0.,0.)	01230
	NSUBS=MINO(NSUBS+1,NPREIM)	01240
	GO TO 530	01250
510	CPRY(IS,NY)=CJ(NPTCHS-NPREJ+J-NSUBS)	01260
530	CONTINUE	01270
	WRITE(6,16)	01280
16	FORMAT('O*****NATURAL MODE*****',/, 'OX-DIRECTED CURRENT')	01290
	DO 540 I=1,NX	01300
	WRITE(6,17) (CPRX(I,J),J=1,NY)	01310
17	FORMAT('O',5(2D12.4,2X))	01320
540	CONTINUE	01330
	WRITE(6,18)	01340
18	FORMAT('OY-DIRECTED CURRENT')	01350
	DO 550 I=1,NX	01360
	WRITE(6,17) (CPRY(I,J),J=1,NY)	01370
550	CONTINUE	01380
	WRITE(6,19)	01390
19	FORMAT('OHOMOGENEOUS EXPANSION COEF''S')	01400
	WRITE(6,17) (CJ(2*NPTCHS-NPRE+I),I=1,NPRE)	01410
	RETURN	01420
	END	01430

C	COMPLEX FUNCTION CPLATE*16(CSUNOR)	01440
C	DETERMINANT EVALUATION ROUTINE FOR HALLEN-TYPE AUGMENTED MOMENT	01450
C	MATRIX FOR THE THIN PLATE SCATTERER	01460
C	FOR S E M APPLICATIONS	01470
C	BY L W PEARSON 8/74	01480
		01490
	IMPLICIT COMPLEX*16(C),REAL*8(A,B,D-H,O-Z)	01500
	COMMON /GEOM/ XSYM,YSYM,W,NX,NY,IPREAS(10),JPREAS(10),NPRED,NPREJ	01510
	COMMON /MAT/ CMATX(25,25),CMATY(25,25),CHOM(50,10),CMAT4(10,25),	01520
	INPTCHS,NDIMI,NDIMCI,NDIMCJ,NORD	01530
	REAL*8 DRARG,DIMARG,DRBES(20),DIMBES(20),DUM1(20),DUM2(20),DUM3(20)	01540
	1),DUM4(20)	01550
	DIMENSION CINTER(10,10),CINTX(25),CINTY(25),CCOSTM(10),CSINTM(10)	01560
	INTEGER MES(4,2)/'SYMM','ETRI','C',' ','ANTI','SYMM','ETRI','C'/	01570
	DATA C/(3.D08,0.D0)/,PLUS/'+'/,PI/3.141592653589793/	01580
	NDIMI=25	01590
	NDIMCI=50	01600
	NDIMCJ=10	01610
	NDIM=50	01620
C		01630
C	FORMULATION SETUP ROUTINES	01640
C		01650
	IMX=1	01660
	IMY=1	01670
	IF(XSYM.NE.PLUS) IMX=2	01680
	IF(YSYM.NE.PLUS) IMY=2	01690
C		01700
C		01710
C		01720
	IM(X/Y)=2 INDICATES ANTISYMMETRIC	01700
	DISTR OF X-DIRECTED CURRENT WRT X	01710
	/Y AXIS	01720
C	NPTCHS=NX*NY	01730
		01740
	TOT NO OF CURRENT PATCHES	01750
		01760
		01770
C		01780
	PROPORTIONAL WIDTH OF EDGE PULSES	01780
C	EDG2=EDGFAC*EDGFAC	01790
		01800
	CORNER FACTOR	01810
		01820
		01830
		01840
		01850
C		01860
	NORMALIZED LAPLACE VARIABLE	01870
		01880
		01890
C		01900
	NUMER INTEG PARMS	01910
		01920
C		01930
	SIGNED SYMMETRY INDICATORS	01940
		01950
		01960
C		01970
C		01980
	SIGNS OF KERNEL FOR EA QUAD'S CON	01970
	TRIBUTION	01980
		01990
		02000
C		02010
	NINDEX = 1 INDICATES EVEN-EVEN OR	02010
	ODD-ODD SYMMETRY FOR X-DIR CURR	02020
		02030
C		02040
		02040
	NSCOS=1	02030
	IF(XSYM.EQ.PLUS) NSCOS=2	02040

C		NSCDS = 2 INDICATES EVEN SYMM WRT	02050
C		Y FOR X DIR CURR (I E COSINE EXPA	02060
C		NSION OF HOMOGENEOUS SOL'N)	02070
C			02080
C			02090
C			02100
C			02110
C			02120
C			02130
C			02140
C			02150
C			02160
C			02170
C			02180
C			02190
C			02200
C			02210
C			02220
C			02230
C			02240
C			02250
C			02260
C			02270
C			02280
C			02290
C			02300
C			02310
C			02320
C			02330
C			02340
C			02350
C			02360
C			02370
C			02380
C			02390
C			02400
C			02410
C			02420
C			02430
C			02440
C			02450
C			02460
C			02470
C			02480
C			02490
C			02500
C			02510
C			02520
C			02530
C			02540
C			02550
C			02560
C			02570
C			02580
C			02590
C			02600
C			02610
C			02620
C			02630
C			02640
C			02650

C	CO=CINTER(ID1,JD1)+NSMIII*CINTER(ID2,JD2)	03270
	SUM OF SOURCE CONT FROM QI & QIII	03280
C	CE=NSMII*CINTER(ID2,JD1)+NSMIV*CINTER(ID1,JD2)	03290
	SUM OF SOURCE CONT FROM QII & QIV	03300
C	CMATY(I,J)=(CO-CE)*EDGFAC	03310
C	SUBMAT ENTRY FOR Y-DIR CURR'S	03320
C	NOTE THAT EVEN Q'S CONT NEGATIVE	03330
	FOR Y-DIR CURR'S	03340
	IF(JS.NE.JPREAS(NSUBS+1)) GO TO 325	03350
	NSUBS=MINO(NSUBS+1,NPREJM)	03360
	GO TO 330	03370
325	CMATX(I,J-NSUBS)=(CE+CO)*EDGFAC	03380
	SUBMAT ENTRY FOR X-DIR CURR'S	03390
C		03400
C	END ROUTINE FOR ABS(X)=A EDGE TERMS	03410
C		03420
330	CONTINUE	03430
C		03440
C	END LOOP OVER Y COORD FOR INTERIOR PATCHES	03450
C		03460
C	BEGIN ROUTINE FOR CONSTRUCTION OF SOURCE TERMS FOR ABS(Y)=B EDGE	03470
C		03480
	JD1=IABS(NY-JM)+1	03490
	JD2=NY+JM	03500
	NSUBSJ=NSUBS	03510
	NSUBS=0	03520
	DO 340 IS=1,NXM1	03530
C		03540
C	INDEX DOWN X COORD'S INTERIOR	03550
	PATCHES	03560
	ID1=IABS(IS-IM)+1	03570
	ID2=IS+IM	03580
	J=(NYM1)*NX+IS	03590
C	CO=CINTER(ID1,JD1)+NSMIII*CINTER(ID2,JD2)	03600
	SUM OF SOURCE CONT FROM QI & QIII	03610
C	CE=NSMII*CINTER(ID2,JD1)+NSMIV*CINTER(ID1,JD2)	03620
	SUM OF SOURCE CONT FROM QII & QIV	03630
C	CMATX(I,J-NSUBSJ)=(CE+CO)*EDGFAC	03640
	SUBMAT ENTRY FOR X-DIR CURR'S	03650
	IF(IS.NE.IPREAS(NSUBS+1)) GO TO 335	03660
	NSUBS=MINO(NSUBS+1,NPREIM)	03670
	GO TO 340	03680
335	CMATY(I,J-NSUBS)=(CO-CE)*EDGFAC	03690
	SUBMAT ENTRY FOR Y-DIR CURR'S	03700
C		03710
C	NOTE THAT EVEN Q'S CONT NEGATIVE	03720
C	FOR Y-DIR CURR'S	03730
340	CONTINUE	03740
C		03750
C	END ROUTINE FOR ABS(Y)=B EDGE	03760
C		03770
C	CONSTRUCTION OF CORNER SOURCE TERM	03780
C		03790
	ID1=IABS(NX-IM)+1	03800
	ID2=NX+IM	03810
	J=NX*NY	03820
	CO=CINTER(ID1,JD1)+NSMIII*CINTER(ID2,JD2)	03830
C		03840
	SUM OF SOURCE CONT FROM QI & QIII	03850
C	CE=NSMII*CINTER(ID2,JD1)+NSMIV*CINTER(ID1,JD2)	03860
	SUM OF SOURCE CONT FROM QII & QIV	03870
C	IF(NY.NF.JPREAS(NPREJ)) CMATX(J,J-NPREJM)=(CE+CO)*EDG2	03880
	SUBMAT ENTRY FOR X-DIR CURR'S	03890
C	IF(NX.NE.IPREAS(NPREI)) CMATY(I,J-NPREIM)=(CO-CE)*EDG2	03900

C		SUBMAT ENTRY FOR Y-DIR CURR'S	03880
C		NOTE THAT EVEN Q'S CONT NEGATIVE	03890
C		FOR Y-DIR CURR'S	03900
C	350	CONTINUE	03910
C		END OF MOMENT MATRIX INTERACTION CONSTRUCTION	03920
C		BEGIN ROUTINE TO ENTER HOMOGENEOUS SOL'N EXPANSION COL'S IN MATRIX	03930
C			03940
C	360	NBES=2*NPRE	03950
C		HIGHEST ORDER BESSEL FUNCTION IN	03960
C		HOMOGENEOUS SOL'N EXPANSION	03970
C		IF(NINDX.EQ.2) NBES=NBES-1	03980
C		ONE LESS IF EVEN INDEX EXPANSION	03990
C		DO 400 IM=1,NX	04000
C		X=(FLOAT(IM)-0.5DO)*DX	04010
C		DO 400 JM=1,NY	04020
C		Y=(FLOAT(JM)-0.5DO)*DY	04030
C		I=(JM-1)*NX+IM	04040
C		INDEXING THRU MATCH-PTS	04050
C		PHI=DATAN(Y/X)	04060
C		RHO=DSQRT(X*X+Y*Y)	04070
C		POLAR COORD'S OF MATCH-PTS	04080
C		DRARG=2*DIMAG(CS)*RHO	04090
C		DIMARG=-2*DREAL(CS)*RHO	04100
C		ARGUMENT OF BESSEL FN'S	04110
C		IF(DABS(DIMARG/DRARG).LT.1.E-20) GO TO 364	04120
C		IF REAL ARG SKIP TO REAL BES CALL	04130
C		CALL BSCJZ(DRARG, DIMARG, DRBES, DIMBES, NBES, 0.DO, 16, IERP, DUM1, DUM2, DUM3, DUM4)	04140
C		GET TABLE OF BESSEL FUNCTIONS	04150
C		GO TO 368	04160
C	364	CALL BSLJZ(DRARG, DRBES, NBES, 0.DO, 16, IERR, DUM1, DUM2)	04170
C		CALL ZEROZ(DIMBES, 2*(NBES+1))	04180
C		CALL ZEROZ(DIMBES, 2*(NBES+1))	04190
C		SET UP PURE REAL BES FUNCTIONS	04200
C	368	CCOSTM(1)=0	04210
C		CSINTM(1)=0	04220
C		ZERO 1ST TERM COEF CONSTRUCTION	04230
C		VECTORS	04240
C		DO 370 II=1,NPREP1	04250
C		INDEX THRU CALC OF COEF CONSTR	04260
C		VECTOR	04270
C		INDX=2*II-NINDX	04280
C		CALC SERIES INDEX	04290
C		IF(INDX.EQ.0) GO TO 370	04300
C		SKIP CALC OF BELOW TERM FOR ZERO	04310
C		INDEX - IT WAS SET TO ZERO ABOVE	04320
C		ARG=DFLOAT(INDX-1)*PHI	04330
C		ARGUMENT OF SIN FN	04340
C		CBES=DCMPLX(DRBES(INDX), DIMBES(INDX))	04350
C		CCOSTM(II)=DCOS(ARG)*CBES*4*PI	04360
C		CSINTM(II)=DSIN(ARG)*CBES*4*PI	04370
C		CALC COEFF CONSTRUCTION TERMS	04380
C	370	CONTINUE	04390
C		DO 380 JJ=1,NPREJ	04400
C		LOOP TO REPLACE COL'S FOR PREASSI	04410
C		GNED J TERMS	04420
C		J=JPREAS(JJ)*NX	04430
C		INDEX OF COL BEING REPLACED	04440
C		INDX=2*JJ-NINDX	04450
C			04460
C			04470
C			04480

C		SERIES INDEX FOR REPLACING TERM	04490
	GO TO (371,372),NSCOS		04500
C		SELECT PROPER SERIES COEF ACCORDI	04510
C		NG TO Y SYMMETRY CONDITION	04520
C		NSCOS=2 INDICATES COSINES IN X	04530
C		CURRENT EQ	04540
371	CHOM(I,JJ)=-PI/(2*CS)*(0.DO,1.DO)**INDX*(CSINTM(JJ)-CSINTM(JJ+1))		04550
	CHOM(NPTCHS+I,JJ)=-PI/(2*CS)*(0.DO,1.DO)**INDX*(CCOSTM(JJ)+		04560
	1CCOSTM(JJ+1))		04570
	GO TO 380		04580
372	CHOM(I,JJ)=-PI/(2*CS)*(0.DO,1.DO)**(INDX+1)*(CCOSTM(JJ+1)-		04590
	1CCOSTM(JJ))		04600
	CHOM(I+NPTCHS,JJ)=-PI/(2*CS)*(0.DO,1.DO)**(INDX+1)*(CSINTM(JJ)+		04610
	1CSINTM(JJ+1))		04620
380	CONTINUE		04630
	DO 390 II=1,NPREI		04640
	J=(NY-1)*NX+IPREAS(II)+NPTCHS		04650
C		LOOP TO REPLACE COL'S FOR PREASSI	04660
	JJ=II+NPRES		04670
C		GNED I TERMS	04680
	INDX=2*(II+NPRES)-NINDX		04690
	GO TO (381,382),NSCOS		04700
381	CHOM(I,JJ)=-PI/(2*CS)*(0.DO,1.DO)**INDX*(CSINTM(II+NPRES)-		04710
	1CSINTM(II+NPRES+1))		04720
	CHOM(I+NPTCHS,JJ)=-PI/(2*CS)*(0.DO,1.DO)**INDX*(CCOSTM(II+NPRES)+		04730
	1CCOSTM(II+NPRES+1))		04740
	GO TO 390		04750
382	CHOM(I,JJ)=-PI/(2*CS)*(0.DO,1.DO)**(INDX+1)*(CCOSTM(II+NPRES+1)-		04760
	1CCOSTM(II+NPRES))		04770
	CHOM(I+NPTCHS,JJ)=-PI/(2*CS)*(0.DO,1.DO)**(INDX+1)*		04780
	1(CSINTM(II+NPRES)+CSINTM(II+NPRES+1))		04790
390	CONTINUE		04800
400	CONTINUE		04810
C			04820
C	END OF MOMENT MATRIX CONSTRUCTION		04830
C			04840
405	CONTINUE		04850
	CALL SPRHOM(CMATX,NPTCHS,NPTCHS-NPREJ,NDIM1,NDIM1,		04860
1	CMATY,NPTCHS,NPTCHS-NPREI,NDIM1,NDIM1,		04870
2	CHOM,NDIMCI,NDIMCJ,CMAT4,NDIMCJ,NDIM1,CDET)		04880
	FRAT=CDABS(CMATX(1,1))		04890
	CPLATE=CDET/FRAT		04900
	WRITE(6,20) CSUNOR,CPLATE		04910
20	FORMAT(' ',5X,2E12.4,5X,2E12.4)		04920
	RETURN		04930
	END		04940

	SUBROUTINE SPRHOM(CMAT1,NI1,NJ1,NDIM1I,NDIM1J,CMAT2,NI2,NJ2,NDIM2I	00010
	1,NDIM2J,CMAT3,NDIM3I,NDIM3J,CMAT4,NDIM4I,NDIM4J,CDET)	00020
	IMPLICIT COMPLEX*16(C),REAL*8(A,B,D-H,C-Z)	00030
C		00040
C	SUBROUTINE TO DIAGCNALIZE AND CALC DETERMINANT OF A SPARCELY-	00050
C	COUPLED MATRIX	00060
C	BY L W PEARSON 7/74	00070
C	REVISED 5/75	00075
C		00080
	DIMENSION CMAT1(NDIM1I,NDIM1J),CMAT2(NDIM2I,NDIM2J),CMAT3(NDIM3I,	00090
	INDIM3J),CMAT4(NDIM4I,NDIM4J)	00100
	NI3=NI1+NI2	00110
	NJ3=NI3-NJ2-NJ1	00120
	CALL ZEROZ(CMAT4,4*NDIM4I*NDIM4J)	00130
	CDET=1	00140
C		00150
	INITIALIZE PRODUCT ACCUMULATOR	
	NPR=3	
	NJ1M1=NJ1-1	00160
	NJ1L=NJ1	
	IF(NJ2+NJ3.GE.1) GO TO 95	
	NJ1L=NJ1L-1	
	NPR=1	
95	DO 155 M=1,NJ1L	00170
C		00180
	INDEX ACROSS COL	
	MP1=M+1	00190
	FMAX=CDABS(CMAT1(M,M))	00200
	K=M	00210
	IF(MP1.GT.NI1) GO TO 105	00220
	DO 100 I=MP1,NI1	00230
C		00240
C		00250
	LOOP TO SEARCH FOR PIVOT IN MTH	
	COL	00260
	FCK=CDABS(CMAT1(I,M))	00270
	IF(FCK.LE.FMAX) GO TO 100	00280
	K=I	00290
C		00300
	IF LARGER ELEMENT FOUND MARK ROW	
	FMAX=FCK	00310
C		00320
C		00330
100	CONTINUE	00340
105	CSTOR=CMAT1(K,M)	00350
C		00360
	CDET=CDET*CSTOR	
C		00370
	MULT PIVOT INTO PROD ACCUMULATOR	
	IF(K.EQ.M) GO TO 115	00380
C		00390
	IF PIVOT ON DIAG SKIP ROW EXCH	
	CDET=-CDET	00400
C		00410
107	DO 110 J=M,NJ1	00420
C		00430
	LOOP TO EXCH DIAG AND PIVOT ROWS	
	CSTO=CMAT1(K,J)	00440
	CMAT1(K,J)=CMAT1(M,J)	00450
	CMAT1(M,J)=CSTO	00460
110	CONTINUE	00470
	IF(NJ3.LT.1) GO TO 115	00475
	DO 112 J=1,NJ3	00480
	CSTO=CMAT3(K,J)	00490
	CMAT3(K,J)=CMAT3(M,J)	00500
	CMAT3(M,J)=CSTO	00510
112	CONTINUE	00520
115	CONTINUE	00560
	IF(MP1.GT.NI1) GO TO 155	00570

C	DO 150 I=MP1,NI1	ELIMINATION LOOP FOR CMAT1	00580
	CFAC=CMAT1(I,M)/CSTOR		00590
C		ELIMINATION FACTOR	00600
	IF(MP1.GT.NJ1) GO TO 125		00610
	DO 120 J=MP1,NJ1		00620
C		LOOP ACROSS ROW IN CMAT3	00630
	CMAT1(I,J)=CMAT1(I,J)-CMAT1(M,J)*CFAC		00640
120	CONTINUE		00650
	IF(NJ3.LT.1) GO TO 150		00660
125	DO 130 J=1,NJ3		00665
C		LOOP ACROSS ROW IN CMAT3	00670
	CMAT3(I,J)=CMAT3(I,J)-CMAT3(M,J)*CFAC		00680
130	CONTINUE		00690
150	CONTINUE		00700
155	CONTINUE		00720
	NI4=NI1-NJ1		00730
	IF(NI4.LE.0) GO TO 290		00740
C			00750
C			00760
C	BEGIN ROUTINE TO CREATE/'DIAGONALIZE' CMAT4		00770
C			00780
	NPIV=NI4		00790
	IF(NI4.GT.NJ2) NPIV=NJ2		00800
	DO 250 M=1,NPIV		00810
C		INDEX ACROSS COL FOR CMAT4 DIAG	00820
	MP1=M+1		00830
	FMAX=CDABS(CMAT2(1,M))		00840
	K=1		00850
	IF(NI2.LT.2) GO TO 205		00860
	DO 200 I=2,NI2		00870
C		LOOP TO SEARCH FOR PIVOT IN MTH	00880
C		COL	00890
	FCK=CDABS(CMAT2(I,M))		00900
	IF(FCK.LE.FMAX) GO TO 200		00910
	K=I		00920
C		IF LARGER ELEMENT FOUND MARK ROW	00930
	FMAX=FCK		00940
C		USE NEW LARGE ELEMENT AS COMPARI-	00950
C		SON VALUE	00960
200	CONTINUE		00970
205	CSTOR=CMAT2(K,M)		00980
C		SAVE VAL OF PIVOT ELEMENT	00990
	CDET=CDET*CSTOR		01000
C		MULT PIVOT INTO PROD ACCUM	01010
	CDET=-CDET		01020
C		CHANGE SIGN OF DETERM BECAUSE OF	01030
C		EXCHANGE FROM CMAT2 TO CMAT4	01040
	DO 210 J=M,NJ2		01050
C		LOOP TO EXCHANGE DIAG AND PIVOT R	01060
C		ROWS	01070
	CSTO=CMAT4(M,J)		01080
	CMAT4(M,J)=CMAT2(K,J)		01090
	CMAT2(K,J)=CSTO		01100
210	CONTINUE		01110
	K3=K+NI1		01120
	M3=NJ1+M		01130
	IF(NJ3.LT.1) GO TO 213		01140
	DO 212 J=1,NJ3		01150
	CSTO=CMAT3(K3,J)		01160
	CMAT3(K3,J)=CMAT3(M3,J)		01170

	CMAT3(M3,J)=CSTO	01180
212	CONTINUE	01190
213	IF(NI2.LT.1) GO TO 290	01225
235	DO 250 I=1,NI2	01230
C		01240
C	LOOP TO CARRY ELIMINATION INTO CMAT2	01250
	I3=NI1+I	01260
	CFAC=CMAT2(I,M)/CSTOR	01270
	IF(MP1.GT.NJ2) GO TO 242	01280
	DO 240 J=MP1,NJ2	01290
C		01300
	LOOP ACROSS ROW OF CMAT2	01310
	CMAT2(I,J)=CMAT2(I,J)-CMAT4(M,J)*CFAC	01320
240	CONTINUE	01325
	IF(NJ3.LT.1) GO TO 250	01330
242	DO 245 J=1,NJ3	01340
C		01350
	LOOP ACROSS ROW OF CMAT3	01360
	CMAT3(I3,J)=CMAT3(I3,J)-CMAT3(M3,J)*CFAC	01380
245	CONTINUE	01390
250	CONTINUE	01400
C		01410
C	END ROUTINE TO 'DIAGONALIZE' CMAT4	01420
C		01430
290	IF(NI4.GE.NJ2) GO TO 390	01440
C		01450
C	IF DIAGONAL DOES NOT PASS THRU SKIP DIAGONALIZATION FOR CMAT2	01460
C		01470
C	BEGIN ROUTINE TO 'DIAGONALIZE' CMAT2	01480
	NI4P1=NI4+1	01482
	NJ2L=NJ2	01484
	IF(NJ3.GE.1) GO TO 295	01486
	NJ2L=NJ2L-1	01488
	NPR=2	01492
295	DO 350 M=NI4P1,NJ2L	01500
	MI=M-NI4	01510
	M3=MI+NI1	01515
	MP1=M+1	01520
	MIP1=MI+1	01530
	FMAX=CDABS(CMAT2(MI,M))	01540
	K=MI	01550
	IF(MIP1.GT.NI2) GO TO 305	01560
	DO 300 I=MIP1,NI2	01570
C		01580
C	LOOP TO SEARCH FOR PIVOT IN MTH COL	01590
	FCK=CDABS(CMAT2(I,M))	01600
	IF(FCK.LE.FMAX) GO TO 300	01610
	K=I	01620
C		01630
	IF LARGER ELEMENT FOUND MARK ROW	01640
C		01650
	USE NEW LARGE ELEMENT AS COMPARI- SON VALUE	01660
300	CONTINUE	01670
305	CSTOR=CMAT2(K,M)	01680
C		01690
	SAVE VAL OF PIVOT ELEMENT	01700
	K3=K+NI1	01710
	CDET=CDET*CSTOR	01720
C		01730
	MULT PIVOT INTO PROD ACCUMULATOR	01740
C		01750
	IF PIVOT ON DIAG SKIP ROW EXCH	
C		
	CDET=-CDET	
C		
	CHANGE SIGN BECAUSE OF ROW EXCH	

	DO 410 J=M,NJ3		02390
C		LOOP TO EXCH DIAG AND PIVOT ROWS	02400
	CSTO=CMAT3(K,J)		02410
	CMAT3(K,J)=CMAT3(MI,J)		02420
	CMAT3(MI,J)=CSTO		02430
410	CONTINUE		02440
415	CONTINUE		02480
	DO 450 I=MIP1,NI3		02490
C		ELIMINATION LOOP	02500
	CFAC=CMAT3(I,M)/CSTOR		02510
	DO 445 J=MP1,NJ3		02520
C		LOOP ACROSS ROW IN CMAT3	02530
	CMAT3(I,J)=CMAT3(I,J)-CMAT3(MI,J)*CFAC		02540
445	CONTINUE		02550
450	CONTINUE		02570
455	GO TO (461,462,463), NPR		02572
461	CDET=CDET*CMAT1(NI1,NJ1)		02574
	RETURN		02576
462	CDET=CDET*CMAT2(NI2,NJ2)		02578
	RETURN		02582
463	CDET=CDET*CMAT3(NI3,NJ3)		02584
	RETURN		02600
C		MULT LAST ELEMENT INTO DETERM	02590
	END		02610

	SUBROUTINE SPRSLV (CMAT1,NI1,NJ1,NDIM1I,NDIM1J,CMAT2,NI2,NJ2,NDIM2I	00010
	1,NDIM2J,CMAT3,NDIM3I,NDIM3J,CMAT4,NDIM4I,NDIM4J,CRHS,CSOLN)	00020
C		00030
C	SUBROUTINE TO BACKSOLVE A TRIANGULARIZED SYSTEM OF SPARCELY-	00040
C	COUPLED LINEAR EQUATION	00050
C	BY L W PEARSON 7/74	00052
C	REVISED 5/75	00054
C		00056
C	STORAGE FORM COMPATIBLE WITH THE TRIANGULARIZATION ROUTINE SPARCE	00060
C		00070
C	THE ENTRY 'HOMSLV' BELOW ALLOWS THE SOLUTION FOR NATURAL VECTORS	00080
C	OF HOMOGENEOUS SYSTEMS PROVIDED THE DETERMINANT OF THE SYSTEM IS	00090
C	ZERO	00100
C		00110
	IMPLICIT COMPLEX*16(C),REAL*8(A,B,D-H,G-Z)	00120
	DIMENSION CMAT1(NDIM1I,NDIM1J),CMAT2(NDIM2I,NDIM2J),CMAT3(NDIM3I,N	00130
	NDIM3J),CMAT4(NDIM4I,NDIM4J),CRHS(NDIM3I),CSOLN(NDIM3I)	00140
	LOGICAL LHOM	00150
C		00160
C	SETUP FOR INHOMOGENECUS SYSTEM	00170
C		00180
	LHOM=.FALSE.	00190
C		00200
	NI3=NI1+NI2	00210
C		00220
	NJ3=NI3-NJ1-NJ2	00230
C		00240
	NI4=NI1-NJ1	00250
C		00260
	ND2=NJ2-NI4	00270
C		00280
	NPR=3	00290
C		00300
	IF(NJ3.LT.1) NPR=2	
C		
	SET INDICATOR FOR NULL CMAT3	
C		
	DEGENERACY	
C		
	IF(NJ3+NJ2.LT.1) NPR=1	
C		
	SET INDICATOR FOR NULL CMAT2 &	
C		
	CMAT3	
C		
	GO TO (81,82,83) , NPR	
C		
	GO MAKE FIRST DIVISION FOR RIGHT-	
C		
	MOST MATRIX	
81	CSOLN(NI3)=CRHS(NI3)/CMAT1(NI1,NJ1)	
	GO TO 100	
82	CSOLN(NI3)=CRHS(NI3)/CMAT2(NI2,NJ2)	
	GO TO 100	
83	CSOLN(NI3)=CRHS(NI3)/CMAT3(NI3,NJ3)	
C		00320
	SOLVE FOR 'LAST' UNKNOWN	
C		00330
	GO TO 100	
C		00340
	GO TO SOLN ROUTINES	
C		00350
	END OF SETUP FOR INHOM SYSTEM	00360
C		00370
	BEGIN ENTRY/SETUP FOR HOMOGENEOUS SYSTEM	00380
C		00390
	ENTRY HOMSLV(CMAT1,NI1,NJ1,NDIM1I,NDIM1J,CMAT2,NI2,NJ2,NDIM2I,NDIM	00400
	12J,CMAT3,NDIM3I,NDIM3J,CMAT4,NDIM4I,NDIM4J,CSOLN,NORD)	00410
C		00420
	LHOM=.TRUE.	00430
C		00440
	LOGICAL INDICATOR FOR HOMOGEN SYS	

	NI3=NI1+NI2	00450
	NJ3=NI3-NJ1-NJ2	00460
	NI4=NI1-NJ1	00470
	ND2=NJ2-NI4	00480
	CSOLN(NI3)=1	00490
C		ASSIGN ARBITRARY ELEMENT IN SOL'N
C		00500
C	END SETUP FOR HOMOGENECUS ENTRY	00510
C		00520
C	BEGIN BACKSOLVE FOR EQUATIONS INVOLVING ONLY CMAT3 (LAST NJ3 EQS)	00530
C		00540
100	FMAX=CDABS(CSOLN(NI3))	00550
	IMAX=NI3	00560
	IF(NJ3.LT.2) GO TO 200	00570
C		00580
C		SKIP ROUTINE IF ONLY LAST VARIABL
C		COUPLES (IT WAS SOLVED/ASSIGNED
		ABOVE)
		00590
	DO 150 IC=2,NJ3	00600
	ICM1=IC-1	00610
	I=NI3-IC+1	00620
	I=NI3-IC+1	00630
		00640
C		00650
C		CALC MATRIX ROW INDX FROM
		COMPLEMENTARY INDX
		00660
	JD3=I-NJ1-NJ2	00670
		00680
C		COL INDX FOR CMAT3 WHICH DEFINES
C		DIAG OF MATRIX
		00690
	CSUM=0	00700
	DO 110 J3C=1,ICM1	00710
		00720
C		LOOP TO ACCUM NEGATIVE SUM OF
C		PREVIOUSLY CALC'D UNKNS
		00730
	J3=NJ3+1-J3C	00740
		00750
C		COL OF COEF IN CMAT3
	J=NI3+1-J3C	00760
		00770
C		ROW OF UNKN IN CSOLN
	CSUM=CSUM-CMAT3(I,J3)*CSOLN(J)	00780
110	CONTINUE	00790
	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I)	00800
C		00810
	CSOLN(I)=CSUM/CMAT3(I,JD3)	ADD R H S TO SUM
		00820
C		DIVIDE BY DIAG COEF
	IF(CDABS(CSOLN(I)).LE.FMAX) GO TO 150	00830
	FMAX=CDABS(CSOLN(I))	00840
	IMAX=I	00850
		00860
C		00870
150	CONTINUE	CHECK FOR MAX ELEMENT
		00880
C		00890
C	BEGIN ROUTINE TO SOLVE FOR ELEMENTS INVOLVING CMAT3 & CMAT2	00900
C		00910
200	IF(NJ3.GE.NI2) GO TO 300	00920
C		00930
C		SKIP ROUTINE IF DIAG DOES NOT
		PASS THRU CMAT2
		00940
	DO 250 IC=1,ND2	00950
	ICM1=IC-1	00960
	I2=NI2-NJ3+1-IC	00970
	I3=NI3-NJ3+1-IC	00980
	JD2=NJ2+1-IC	00990
	NCM1=NJ3+IC-1	01000
	CSUM=0	01010
	IF(NJ3.LT.1) GO TO 215	01020
	DO 210 JC=1,NJ3	
C		01030
		LOOP TO SUM CONTRIB FROM CMAT3
		01040

	J3=NJ3+1-JC	01050
	J=NI3+1-JC	01060
	CSUM=CSUM-CMAT3(I3,J3)*CSOLN(J)	01070
210	CONTINUE	01080
215	IF(ICM1.LT.1) GO TO 225	01090
C	SKIP IF NO TERMS CONTRIB FR CMAT2	01100
	DO 220 J2C=1,ICM1	01110
	J2=NJ2+1-J2C	01120
	J=NI3-NJ3+1-J2C	01130
	CSUM=CSUM-CMAT2(I2,J2)*CSOLN(J)	01140
220	CONTINUE	01150
225	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I3)	01160
	CSOLN(I3)=CSUM/CMAT2(I2,JD2)	01170
	IF(CDABS(CSOLN(I3)).LE.FMAX) GO TO 250	01180
	FMAX=CDABS(CSOLN(I3))	01190
	IMAX=I	01200
250	CONTINUE	01210
C		01220
C	BEGIN ROUTINE TO SOLVE FOR ELEMENTS INVOLVING CMAT3 & CMAT4	01230
C		01240
300	IF(NI4.LT.1) GO TO 400	
	DO 350 IC=1,NI4	01260
	I4=NI4+1-IC	01270
	JD4=I4	01280
	I3=NI1+1-IC	01290
	CSUM=0	
	IF(NJ3.LT.1) GO TO 315	01300
	DO 310 J3C=1,NJ3	01310
	J3=NJ3+1-J3C	01320
	J=NI3+1-J3C	01330
	CSUM=CSUM-CMAT3(I3,J3)*CSOLN(J)	01340
310	CONTINUE	01350
315	NSUBS=ND2+IC-1	01360
C	NO OF NON-DIAG CMAT4 EL'S IN EQ	01370
	IF(NSUBS.LT.1) GO TO 325	01380
	DO 320 J4C=1,NSUBS	01390
	J4=NJ2+1-J4C	01400
	J=NI3-NJ3+1-J4C	01410
	CSUM=CSUM-CMAT4(I4,J4)*CSOLN(J)	01420
320	CONTINUE	01430
325	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I3)	01440
	CSOLN(I3)=CSUM/CMAT4(I4,I4)	01450
	IF(CDABS(CSOLN(I3)).LE.FMAX) GO TO 350	01460
	FMAX=CDABS(CSOLN(I3))	01470
	IMAX=I3	01480
350	CONTINUE	01490
C		01500
C	BEGIN ROUTINE TO SOLVE EQ'S INVOLVING CMAT3 & CMAT1	01510
C		
400	IF(NJ1.LT.1) GO TO 455	01520
	DO 450 IC=1,NJ1	01530
	I=NJ1+1-IC	01540
	ICM1=IC-1	01550
	CSUM=0	
	IF(NJ3.LT.1) GO TO 415	01560
	DO 410 J3C=1,NJ3	01570
	J3=NJ3+1-J3C	01580
	J=NI3+1-J3C	01590
	CSUM=CSUM-CMAT3(I,J3)*CSOLN(J)	01600
410	CONTINUE	01610
415	IF(ICM1.LT.1) GO TO 425	

	DO 420 JC=1, ICM1	01620
	J=NJ1+1-JC	01630
	CSUM=CSUM-CMAT1(I,J)*CSOLN(J)	01640
420	CONTINUE	01650
425	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I)	01660
	CSOLN(I)=CSUM/CMAT1(I,I)	01670
	IF(CDABS(CSOLN(I)).LE.FMAX) GO TO 450	01680
	FMAX=CDABS(CSCLN(I))	01690
	IMAX=I	01700
450	CONTINUE	01710
C		01720
C		01730
C	END OF SOLUTION	01740
C		01750
455	IF(.NOT.LHOM) RETURN	01760
C		01770
C	RETURN IF INHOM SYSTEM	01780
C		01790
C	BEGIN NORMALIZATION ROUTINE FOR NATURAL VECTOR FOR HOMOGENEOUS	01800
C	CASE	01810
	CSCALE=1./CSOLN(IMAX)	01820
	DO 500 I=1, NI3	01830
	CSOLN(I)=CSOLN(I)*CSCALE	01840
500	CONTINUE	01850
	RETURN	01860
	END	01870

```
SUBROUTINE COPYZ(X,Y,N)
DIMENSION X(1),Y(1)
DO 100 I=1,N
X(I)=Y(I)
100 CONTINUE
RETURN
END
```

09340
09350
09360
09370
09380
09390
09400

```
      SUBROUTINE ZEROZ(IARRAY,N)
      DIMENSION IARRAY(1)
      DO 100 I=1,N
      IARRAY(I)=0
100  CONTINUE
      RETURN
      END
```

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09410
09420
09430
09440
09450
09460
09470
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	SUBROUTINE DWEDDL(FCN,N,DELTA,VINT)	09480
	IMPLICIT REAL*8(A-H,O-Z)	09490
	COMPLEX*16 FCN,C,VINT	09500
	DIMENSION FCN(N)	09510
	DIMENSION COEF(6)	09520
	DATA COEF/2.00,5.00,1.00,6.00,1.00,5.00/	09530
	IF((N-1)/6*6.EQ.N-1) GO TO 100	09540
	WRITE(6,1)	09550
1	FORMAT('OINCORRECT POINTS TO WEDDLE')	09560
	A=1/0	09570
100	CONTINUE	09580
	VINT=0	09590
	DO 200 J=1,N	09600
	JCOEF=J-((J-1)/6)*6	09610
	VINT=VINT+COEF(JCOEF)*FCN(J)	09620
200	CONTINUE	09630
	VINT=(VINT-FCN(1)-FCN(N))*(0.300,0.00)*DCMPLX(DELTA,0.00)	09640
	RETURN	09650
	END	09660


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C.ZANLYT.....D.....ZAN09670
C
C FUNCTION - DETERMINATION OF ZEROS OF AN ANALYTIC COMPLEX ZAN09680
C FUNCTION USING MULLER'S METHOD WITH ZAN09690
C DEFLATION ZAN09700
C USAGE - CALL ZANLYT (F, EPS, NSIG, KN, NGUESS, N, X, ITMAX, ZAN09710
C INFER, IER) ZAN09720
C PARAMETERS F - A FUNCTION SUBPROGRAM, F(Z), WRITTEN BY THE ZAN09730
C USER SPECIFYING THE EQUATION WHOSE ROOTS ZAN09740
C ARE TO BE FOUND. F MUST BE TYPE-NAMED AS ZAN09750
C FOLLOWS - COMPLEX FUNCTION F*16 (Z) ZAN09760
C EPS - 1ST STOPPING CRITERION. A ROOT Z IS ACCEPTED ZAN09770
C IF ABSOLUTE VALUE OF F(Z) .LE. EPS (INPUT) ZAN09780
C NSIG - 2ND STOPPING CRITERION. A ROOT IS ACCEPTED ZAN09790
C IF TWO SUCCESSIVE APPROXIMATIONS TO A GIVEN ZAN09800
C ROOT AGREE IN THE FIRST NSIG DIGITS. (INPUT) ZAN09810
C NOTE. IF EITHER OR BOTH OF THE STOPPING ZAN09820
C CRITERIA ARE FULFILLED, THE ROOT IS ZAN09830
C ACCEPTED. ZAN09840
C KN - THE NUMBER OF KNOWN ROOTS WHICH MUST BE STORED ZAN09850
C IN X(1),...,X(KN), PRIOR TO ENTRY TO ZANLYT ZAN09860
C NGUESS - THE NUMBER OF INITIAL GUESSES PROVIDED. THESE ZAN09870
C GUESSES MUST BE STORED IN X(KN+1),..., ZAN09880
C X(KN+NGUESS) AND NGUESS MUST BE SET EQUAL ZAN09890
C TO ZERO IF NO GUESSES ARE PROVIDED. (INPUT) ZAN09900
C N - THE NUMBER OF NEW ROOTS TO BE FOUND BY ZAN09910
C ZANLYT (INPUT) ZAN09920
C X - A LONG-WORD COMPLEX VECTOR ARRAY OF LENGTH ZAN09930
C .GE. 3*(KN+N). X(1),...,X(KN) ON INPUT ZAN09940
C MUST CONTAIN ANY KNOWN ROOTS. X(KN+1),..., ZAN09950
C X(KN+N) ON INPUT MAY, AT THE USER'S OPTION, ZAN09960
C CONTAIN INITIAL GUESSES FOR THE N NEW ZAN09970
C ROOTS WHICH ARE TO BE COMPUTED. ON OUTPUT, ZAN09980
C X(KN+1),..., X(KN+N) CONTAIN EITHER A ROOT ZAN09990
C CORRECT TO WITHIN A CONVERGENCE CRITERION ZAN10000
C OR THE VALUE(12345678.12345678D+0,12345678. ZAN10010
C 12345678D+0) INDICATIVE OF A FAILURE TO ZAN10020
C ACHIEVE THE SPECIFIED CONVERGENCE FOR THAT ZAN10030
C ROOT, SAY X(KN+J). IN THE LATTER CASE, THE ZAN10040
C MOST RECENT APPROXIMATION TO X(KN+J) IS ZAN10050
C AVAILABLE IN X(ISUB), WHERE ISUB=2*(KN+N)+J ZAN10060
C ITMAX - THE MAXIMUM ALLOWABLE NUMBER OF ITERATIONS ZAN10070
C PER ROOT (INPUT) ZAN10080
C INFER - AN INTEGER VECTOR OF LENGTH .GE. KN+N. ON ZAN10090
C OUTPUT INFER(J) CONTAINS THE NUMBER OF ZAN10100
C ITERATIONS USED IN FINDING THE J-TH ROOT ZAN10110
C WHEN CONVERGENCE WAS ACHIEVED. IF ZAN10120
C CONVERGENCE WAS NOT OBTAINED IN ITMAX ZAN10130
C ITERATIONS, INFER(J) WILL CONTAIN ITMAX+1 ZAN10140
C (OUTPUT) ZAN10150
C IER - ERROR PARAMETER (OUTPUT) ZAN10160
C WARNING ERROR = 32 + N ZAN10170
C N = 1 FAILURE TO CONVERGE WITHIN ITMAX ZAN10180
C ITERATIONS FOR ONE OF THE (N) NEW ROOTS TO ZAN10190
C BE FOUND ZAN10200
C PRECISION - DOUBLE ZAN10210
C REQ'D IMSL ROUTINES - UERTST ZAN10220
C AUTHOR/IMPLEMENTOR - C. G. JOHNSON/L. L. WILLIAMS ZAN10230
C LANGUAGE - FORTRAN ZAN10240
C ZAN10250
C .....ZAN10260
C LATEST REVISION - SEPTEMBER 1, 1971 ZAN10270

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C	SUBROUTINE ZANLYT (F, EPS, NSIG, KN, NGUESS, N, X, ITMAX, INFER, IER) COMPLEX*16 1 DOUBLE PRECISION DIMENSION IER = 0 ONE = (1.0D+00, 0.0D+00) EPS1 = 10.0D+00**(-NSIG) ICDNJ = 0 IBOMB = 0	ZAN10280 ZAN10290 ZAN10300 ZAN10310 ZAN10320 ZAN10330 ZAN10340 ZAN10350 ZAN10360 ZAN10370 ZAN10380 ZAN10390 ZAN10400 ZAN10410 ZAN10420 ZAN10430 ZAN10440 ZAN10450 ZAN10460 ZAN10470 ZAN10480 ZAN10490 ZAN10500 ZAN10510 ZAN10520 ZAN10530 ZAN10540 ZAN10550 ZAN10560 ZAN10570 ZAN10580 ZAN10590 ZAN10600 ZAN10610 ZAN10620 ZAN10630 ZAN10640 ZAN10650 ZAN10660 ZAN10670 ZAN10680 ZAN10690 ZAN10700 ZAN10710 ZAN10720 ZAN10730 ZAN10740 ZAN10750 ZAN10760 ZAN10770 ZAN10780 ZAN10790 ZAN10800 ZAN10810 ZAN10820 ZAN10830 ZAN10840 ZAN10850 ZAN10860 ZAN10870 ZAN10880
	SET NUMBER OF ITERATIONS	
C	MB1 = KN+1 MB2 = KN+N LSTART = MB2+1 MPG = MB1+NGUESS DO 2 I = MPG, MB2 2 X(I) = (0.0D+0, 0.0D+0) L = MB1 IF (KN .EQ. 0) GO TO 5 DO 3 I = 1, KN INFER(I) = 0 ITEMP = MB2+I X(ITEMP) = X(I) ITEMP = MB2+ITEMP 3 X(ITEMP) = X(I) 5 JK = 0 OZ = CDABS(X(L)) IF (OZ .LE. 1.0D-15) GO TO 25	
	ROOT ESTIMATE NOT EQUAL TO ZERO	
C	10 RT = (.9D+00, 0.0D+00)*X(L) ASSIGN 15 TO NN GO TO 135 15 XO = FPRT RT = (1.1D+00, 0.0D+00)*X(L) ASSIGN 20 TO NN GO TO 135 20 X1 = FPRT H = X(L)-RT RT = X(L) ASSIGN 40 TO NN GO TO 135	
	ROOT ESTIMATE EQUAL TO ZERO	
C	25 RT = -ONE ASSIGN 30 TO NN GO TO 135 30 XO = FPRT RT = ONE ASSIGN 35 TO NN GO TO 135 35 X1 = FPRT RT = (0.0D+00, 0.0D+00) H = -ONE ASSIGN 40 TO NN GO TO 135 40 X2 = FPRT 45 D = (-0.5D+00, 0.0D+00)	
	BEGIN MAIN ALGORITHM	
C	50 DD = ONE + D T1 = XO*D*DD T2 = X1*DD*DD	

	XX = X2*DD	ZAN10890
	T3 = X2*D	ZAN10900
	BI = T1-T2+XX+T3	ZAN10910
	DEN = BI*BI-(4.0D+0,0.0D+0)*(XX*T1-T3*(T2-XX))	ZAN10920
C		USE DENOMINATOR OF MAXIMUM AMPLITUDE ZAN10930
	T1 = CDSQRT(DEN)	ZAN10940
	T2 = BI + T1	ZAN10950
	T3 = BI - T1	ZAN10960
	QZ = CDABS(T2) - CDABS(T3)	ZAN10970
	IF (QZ .GE. 0) GO TO 60	ZAN10980
55	DEN = T3	ZAN10990
	GO TO 65	ZAN11000
60	DEN = T2	ZAN11010
		TEST FOR ZERO DENOMINATOR ZAN11020
65	QZ = CDABS(DEN)	ZAN11030
	IF (QZ .GT. 1.0D-15) GO TO 75	ZAN11040
70	DEN = ONE	ZAN11050
75	DI = ((-2.0D+00,0.0D+00)*XX)/DEN	ZAN11060
	H = DI * H	ZAN11070
	RT = RT + H	ZAN11080
		CHECK CONVERGENCE OF THE FIRST KIND ZAN11090
	QZ = CDABS(H/RT)	ZAN11100
	IF (QZ .LE. EPS1) GO TO 100	ZAN11110
80	ASSIGN 85 TO MN	ZAN11120
	GO TO 135	ZAN11130
85	QZ = CDABS(FPRT)-CDABS(X2*(10.0D0,0.0D0))	ZAN11140
	IF (QZ .LT. 0.0D+0) GO TO 95	ZAN11150
		TAKE REMEDIAL ACTION TO INDUCE ZAN11160
		CONVERGENCE ZAN11170
90	DI = DI*(0.5D+00,0.0D+00)	ZAN11180
	H = H*(0.5D+00,0.0D+00)	ZAN11190
	RT = RT-H	ZAN11200
	GO TO 135	ZAN11210
95	X0 = X1	ZAN11220
	X1 = X2	ZAN11230
	X2 = FPRT	ZAN11240
	D = DI	ZAN11250
	GO TO 50	ZAN11260
		A ROOT HAS BEEN FOUND ZAN11270
100	FRT = F(RT)	ZAN11280
105	X(L) = RT	ZAN11290
	ITEMP = MB2+L-IBOMB	ZAN11300
	X(ITEMP) = RT	ZAN11310
	ITEMP = MB2+MB2+L	ZAN11320
	X(ITEMP) = RT	ZAN11330
		CHECK TO SEE IF COMPLEX-CONJUGATE ZAN11340
		IS ALSO A ROOT ZAN11350
	IF (CDABS(F(DCONJG(X(L)))) .GT. 10.0D+0*CDABS(FRT)) GO TO 115	ZAN11360
	QZ = CDABS(X(L)- DCONJG(X(L)))	ZAN11370
	IF (ICONJ .NE. 0 .OR. QZ .LT. 1.0D-8) GO TO 115	ZAN11380
	ISTART = L+2	ZAN11390
	INSER1 = L+1	ZAN11400
	DO 110 INSERT = ISTART,MB2	ZAN11410
	X(INSERT) = X(INSER1)	ZAN11420
110	INSER1 = INSERT	ZAN11430
	X(L+1) = DCONJG(X(L))	ZAN11440
	ICONJ = 1	ZAN11450
	GO TO 120	ZAN11460
115	ICONJ = 0	ZAN11470
120	CONTINUE	ZAN11480
125	INFER(L) = JK	ZAN11490

	L = L+1	ZAN1150C
	IF (L .LE. MB2) GO TO 5	ZAN11510
C	RETURN TO CALLING PROGRAM	ZAN11520
	130 GO TO 185	ZAN11530
	135 JK = JK+1	ZAN11540
	IF (JK .GT. ITMAX) GO TO 180	ZAN11550
	140 FRT = F(RT)	ZAN11560
	FPRT = FRT	ZAN11570
C	TEST TO SEE IF FIRST ROOT IS BEING	ZAN11580
C	DETERMINED	ZAN11590
	IF (L .EQ. 1) GO TO 160	ZAN11600
	IF (L .LE. IBOMB+1) GO TO 160	ZAN11610
C	COMPUTE DENOMINATOR FOR MODIFIED	ZAN11620
C	FUNCTION	ZAN11630
	145 LIMUP = MB2+L-IBOMB-1	ZAN11640
	DJ 150 I = LSTART,LIMUP	ZAN11650
	TEM = RT - X(I)	ZAN11660
	OZ = CDABS(TEM)	ZAN11670
	IF (OZ .LT. 5.00-15) GO TO 175	ZAN11680
	150 FPRT = FPRT/TEM	ZAN11690
C	CHECK CONVERGENCE OF THE SECOND KIND	ZAN11700
	160 OZ = CDABS(FRT)	ZAN11710
	IF (OZ .GE. EPS) GO TO 170	ZAN11720
	165 OZ = CDABS(FPRT)	ZAN11730
	IF (OZ .LT. EPS) GO TO 105	ZAN11740
	170 GO TO NN,(15,20,30,35,40,85)	ZAN11750
	175 RT = RT * (1.000001D+0,0.00+0)	ZAN11760
	GO TO 135	ZAN11770
C	WARNING ERROR, ITMAX = MAXIMUM	ZAN11780
	180 IER = 33	ZAN11790
	INFER(L) = ITMAX + 1	ZAN11800
	IBOMB = IBOMB + 1	ZAN11810
	X(L) = (12345678.12345678D+0,12345678.12345678D+0)	ZAN11820
	ITEMP = MB2 ÷ MB2 + L	ZAN11830
	X(ITEMP) = RT	ZAN11840
	L = L+1	ZAN11850
	IF (L .LE. MB2) GO TO 5	ZAN11860
	185 IF (IER .EQ. 0) GO TO 9005	ZAN11870
	9000 CONTINUE	ZAN11880
	CALL UERTST(IER,'ZANLYT')	ZAN11890
	9005 RETURN	ZAN11900
	END	ZAN11910

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C.UERTST.....UER11920
C
C FUNCTION - ERROR MESSAGE GENERATION UER11930
C USAGE - CALL UERTST(IER,'NAMEXX') UER11940
C PARAMETERS IER - ERROR PARAMETER. TYPE + N WHERE UER11950
C TYPE= 128 IMPLIES TERMINAL ERROR UER11960
C 64 IMPLIES WARNING WITH FIX JER11970
C 32 IMPLIES WARNING UER11980
C N = ERROR CODE RELEVANT TO CALLING ROUTINE UER11990
C NAMEXX - NAME OF THE CALLING ROUTINE JER12000
C AUTHOR/IMPLEMENTER - PEDER SVENDSEN UER12010
C LANGUAGE - FORTRAN UER12020
C.....UER12030
C LATEST REVISION - JANUARY 19, 1971 UER12040
C UER12050
C SUBROUTINE UERTST(IER,NAME) UER12060
C UER12070
C DIMENSION ITYP(5,4),IBIT(4) JER12080
C INTEGER*2 NAME(3) UER12090
C INTEGER WARN,WARF,TERM,PRINTR UER12100
C EQUIVALENCE (IBIT(1),WARN),(IBIT(2),WARF),(IBIT(3),TERM) UER12110
C DATA ITYP /'WARN','ING ',' ',' ',' ',' ' UER12120
C * 'WARN','ING(','','WITH',' FIX','') ' UER12130
C * 'TERM','INAL',' ',' ',' ',' ' UER12140
C * 'NON-','DEFI','NED ',' ',' ',' ' UER12150
C * DATA IBIT / 32,64,128,0/ UER12160
C DATA PRINTR / 6/ UER12170
C IER2=IER UER12180
C IF (IER2 .GE. WARN) GO TO 5 UER12190
C UER12200
C IER1=4 NON-DEFINED UER12210
C GO TO 20 UER12220
C 5 IF (IER2 .LT. TERM) GO TO 10 UER12230
C UER12240
C IER1=3 TERMINAL UER12250
C GO TO 20 UER12260
C 10 IF (IER2 .LT. WARF) GO TO 15 UER12270
C UER12280
C IER1=2 WARNING(WITH FIX) UER12290
C GO TO 20 UER12300
C UER12310
C 15 IER1=1 WARNING UER12320
C UER12330
C 20 IER2=IER2-IBIT(IER1) EXTRACT 'N' UER12340
C UER12350
C PRINT ERROR MESSAGE JER12360
C WRITE (PRINTR,25) (ITYP(I,IER1),I=1,5),NAME,IER2 UER12370
C 25 FORMAT(' *** I M S L(UERTST) *** ',5A4,4X,3A2,4X,12) UER12380
C RETURN UER12390
C END UER12400

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	SUBROUTINE BSLJZ(X,FJ,NMAX,A,ND,IERR,FJAPRX,RR)	00012410
	IMPLICIT REAL*8 (A-H,O-Z)	00012420
	DIMENSION FJ(1),FJAPRX(1),RR(1)	00012430
	NMAXT=NMAX	00012440
	IF(NMAXT.GE.0)GO TO 30	00012450
	IF(DABS(A).LE.1.0D-15)GO TO 10	00012460
	GO TO 20	00012470
10	IERR=4	00012480
	RETURN	00012490
20	NMAXAR=IABS(NMAXT)	00012500
	NMAXT=1	00012510
30	IF(A.GT.0.0)GO TO 40	00012520
	IF(DABS(A).LE.1.0D-15)GO TO 40	00012530
	IERR=1	00012540
	RETURN	00012550
40	IF(A.LT.1.0)GO TO 70	00012560
	IERR=2	00012570
	RETURN	00012580
70	IF(X.GT.0.0)GO TO 130	00012590
	IERR=3	00012600
	RETURN	00012610
130	IERR=0	00012620
	EPSLON=.500*10.**(-ND)	00012630
	NMPI=NMAX+1	00012640
	DO 160 N=1,NMPI	00012650
160	FJAPRX(N)=0.0	00012660
	SUM=(X/2.)**A/DGAMMA(1.+A)	00012670
	D1=2.302600*ND+1.386300	00012680
	IF(NMAXT.LE.0)GO TO 230	00012690
	Y=.500*D1/NMAXT	00012700
	CALL TZ(Y,TANS)	00012710
	R=NMAXT*TANS	00012720
	GO TO 240	00012730
230	R=0.0	00012740
240	Y=.7357600*D1/X	00012750
	CALL TZ(Y,TANS)	00012760
	S=1.359100*X*TANS	00012770
	IF(R.GT.S)GO TO 280	00012780
	NU=1+IDINT(S)	00012790
	GO TO 290	00012800
280	NU=1+IDINT(R)	00012810
290	M=0	00012820
	FL=1.	00012830
	LIMIT=(NU/2)	00012840
320	M=M+1	00012850
	FL=FL*(M+A)/(M+1.00)	00012860
	IF(M.LT.LIMIT)GO TO 320	00012870
	N=2*M	00012880
	R=0.0	00012890
	S=0.0	00012900
390	DENOM=2.*(A+N)/X-R	00012910
	IF(DABS(DENOM).LE.1.0D-15)DENOM=DENOM+1.0D-15	00012920
430	R=1./DENOM	00012930
	NMOD2=MOD(N,2)	00012940
	IF(NMOD2.NE.0)GO TO 480	00012950
	FL=FL*(N+2.00)/(N+2.*A)	00012960
	FLMBDA=FL*(N+A)	00012970
	GO TO 490	00012980
480	FLMBDA=0.0	00012990
490	S=R*(FLMBDA+S)	00013000
	IF(N.LE.NMAXT)RR(N)=P	00013010

	N=N-1	00013020
	IF(N.GE.1)GO TO 390	00013030
	FJ(1)=SUM/(1.+S)	00013040
	IF(NMAXT.EQ.0)GO TO 570	00013050
	DO 560 N=1,NMAXT	00013060
560	FJ(N+1)=RR(N)*FJ(N)	00013070
570	DO 640 N=1,NMP1	00013080
	IF(DABS((FJ(N)-FJAPRX(N))/FJ(N)).LE.EPSLON)GO TO 640	00013090
	DO 610 M=1,NMP1	00013100
610	FJAPRX(M)=FJ(M)	00013110
	NU=NU+5	00013120
	GO TO 290	00013130
640	CONTINUE	00013140
	IF(NMAX.GE.0)RETURN	00013150
	FJ(2)=2.*A*FJ(1)/X-FJ(2)	00013160
	IF(NMAXAB.EQ.1)RETURN	00013170
	DO 650 N=2,NMAXAB	00013180
650	FJ(N+1)=2.*(A-N)*FJ(N)/X-FJ(N-1)	00013190
	RETURN	00013200
	END	00013210

	SUBROUTINE BSCJZ(X,Y,U,V,NMAX,A,ND,IERR,UAPPRX,VAPPRX,RR1,RR2)	BSC13220
	IMPLICIT REAL*8 (A-H,O-Z)	BSC13230
	DIMENSION U(100),V(100),UAPPRX(100),VAPPRX(100),RR1(100),	BSC13240
1	RR2(100)	BSC13250
	IF(A.GE.0.0)GO TO 40	BSC13260
	IERR=1	BSC13270
	RETURN	BSC13280
40	IF(A.LT.1)GO TO 70	BSC13290
	IERR=2	BSC13300
	RETURN	BSC13310
70	IF(X.GT.0.0)GO TO 110	BSC13320
	IF(DABS(Y).LE.1.0D-14)GO TO 90	BSC13330
	GO TO 110	BSC13340
90	IERR=3	BSC13350
	RETURN	BSC13360
110	IF(NMAX.GE.0)GO TO 140	BSC13370
	IERR=4	BSC13380
	RETURN	BSC13390
140	IERR=0	BSC13400
	EPSLON=.5D0*10.**(-ND)	BSC13410
	NMPI=NMAX+1	BSC13420
	DO 200 N=1,NMPI	BSC13430
	UAPPRX(N)=0.0	BSC13440
200	VAPPRX(N)=0.0	BSC13450
	Y1=DABS(Y)	BSC13460
	RZ2=X**2+Y**2	BSC13470
	RZ=DSQRT(RZ2)	BSC13480
	IF(DABS(X).LE.1.0D-14)GO TO 290	BSC13490
	PHI=DATAN2(Y1,X)	BSC13500
	IF(X.LT.0.0) PHI=3.141592653589793D0 + PHI	BSC13510
	GO TO 300	BSC13520
290	PHI=1.570796326794896D0	BSC13530
300	C=DEXP(Y1)*(RZ/2.)**A/DGAMMA(1.+A)	BSC13540
	SUM2=A*PHI-X	BSC13550
	SUM1=C*DCOS(SUM2)	BSC13560
	SUM2=C*DSIN(SUM2)	BSC13570
	D1=2.3026D0*ND+1.3863D0	BSC13580
	IF(NMAX.GT.0)GO TO 380	BSC13590
	R=0.0	BSC13600
	GO TO 390	BSC13610
380	PARAM=.5D0*D1/NMAX	BSC13620
	CALL TZ(PARAM,TANS)	BSC13630
	R=NMAX*TANS	BSC13640
390	S=1.3591D0*RZ	BSC13650
	PARAM=.73576D0*(D1-Y1)/RZ	BSC13660
	CALL TZ(PARAM,TANS)	BSC13670
	IF(Y1.LT.D1)S=S*TANS	BSC13680
	IF(R.GT.S)GO TO 450	BSC13690
	NU=1+IDINT(S)	BSC13700
	GO TO 460	BSC13710
450	NU=1+IDINT(R)	BSC13720
460	N=0	BSC13730
	FL=1.	BSC13740
	C1=1.	BSC13750
	C2=0.	BSC13760
500	V=V+1	BSC13770
	FL=FL*(N+2.*A)/(N+1.D0)	BSC13780
	C=-C1	BSC13790
	C1=C2	BSC13800
	C2=C	BSC13810
	IF(N.LT.NU)GO TO 500	BSC13820

	R1=0.0	BSC13830
	R2=0.0	BSC13840
	S1=0.0	BSC13850
	S2=0.0	BSC13860
610	C=(2.*(A+N)-X*R1+Y1*R2)**2+(X*R2+Y1*R1)**2	BSC13870
	R1=(2.*(A+N)*X-RZ2*R1)/C	BSC13880
	R2=(2.*(A+N)*Y1+RZ2*R2)/C	BSC13890
	FL=FL*(N+1.D0)/(N+2.*A)	BSC13900
	C=2.*(N+A)*FL	BSC13910
	FLAMB1=C*C1	BSC13920
	FLAMB2=C*C2	BSC13930
	C=C1	BSC13940
	C1=-C2	BSC13950
	C2=C	BSC13960
	S=R1*(FLAMB1+S1)-R2*(FLAMB2+S2)	BSC13970
	S2=R1*(FLAMB2+S2)+P2*(FLAMB1+S1)	BSC13980
	S1=S	BSC13990
	IF(N.GT.NMAX)GO TO 770	BSC14000
	RR1(N)=R1	BSC14010
	RR2(N)=R2	BSC14020
770	N=N-1	BSC14030
	IF(N.GE.1)GO TO 610	BSC14040
	C=(1.+S1)**2+S2**2	BSC14050
	U(1)=(SUM1*(1.+S1)+SUM2*S2)/C	BSC14060
	V(1)=(SUM2*(1.+S1)-SUM1*S2)/C	BSC14070
	IF(NMAX.EQ.0)GO TO 850	BSC14080
	DO 840 N=1,NMAX	BSC14090
	U(N+1)=RR1(N)*U(N)-RR2(N)*V(N)	BSC14100
840	V(N+1)=RR1(N)*V(N)+RR2(N)*U(N)	BSC14110
850	IF(Y.LT.0.0)GO TO 860	BSC14120
	GO TO 880	BSC14130
860	DO 870 N=1,NMPL	BSC14140
870	V(N)=-V(N)	BSC14150
880	DO 950 N=1,NMPL	BSC14160
	TEMP1=(U(N)-UAPPRX(N))**2	BSC14170
	TEMP1=TEMP1+(V(N)-VAPPRX(N))**2	BSC14180
	TEMP1=TEMP1/(U(N)**2+V(N)**2)	BSC14190
	IF(TEMP1.LE.EPSLON)GO TO 950	BSC14200
	DO 920 M=1,NMPL	BSC14210
	UAPPRX(M)=U(M)	BSC14220
920	VAPPRX(M)=V(M)	BSC14230
	NU=NU+5	BSC14240
	GO TO 460	BSC14250
950	CONTINUE	BSC14260
	RETURN	BSC14270
	END	BSC14280
	SUBROUTINE TZ(Y,TANS)	BSC14290
	REAL*8 Y,Z,P,TANS,DLOG	BSC14300
	IF(Y.GT.10.0)GO TO 40	BSC14310
	P=.000057941D0*Y-.00176148D0	BSC14320
	P=Y*P+.0208645D0	BSC14330
	P=Y*P-.129013D0	BSC14340
	P=Y*P+.85777D0	BSC14350
	TANS=Y*P+1.10125D0	BSC14360
	RETURN	BSC14370
40	Z=DLOG(Y)-.775D0	BSC14380
	P=(.775D0-DLOG(Z))/(1.0+Z)	BSC14390
	TANS=Y/((1.+P)*Z)	BSC14400
	RETURN	BSC14410
	END	BSC14420

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