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Magnetic Field on a Cylinder in a Loop

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Abstract

The magnetostatic field at the surface of an infinitely long, perfectly conducting, circular cylinder is computed. The source of the field is a circular, current-carrying loop coaxial with the cylinder.

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I. Introduction

An EMP simulator in the shape of a half toroid has been described in Sensor and Simulation Note 94 [Ref. 1]. The low-frequency magnetic field distribution for such a simulator was calculated in Sensor and Simulation Note 112 [Ref. 2]. The effect of a hemispherical test body on the lowfrequency magnetic field was determined in Sensor and Simulation Note 120 [Ref. 3]. In this brief note we will present data on the low-frequency magnetic field when the test body may be idealized as an infinite, perfectly conducting, semicircular cylinder coaxial with the half toroid, the half toroid being in the vertical position.

In making our calculations we will make the same assumptions concerning the effect of the ground conductivity on the magnetic field that were made in Notes 112 and 120. That is to say, firstly, we will assume that the frequency is not exactly zero (which would make it necessary to determine separately the effect of the current distribution in the lower half space on the magnetic field. See, for example, Ref. 4). Secondly, we will assume the frequency is low enough so that, as far as the magnetic field distribution above the ground surface is concerned, the ground may be thought of as perfectly conducting and image theory may be applied. The above two assumptions are compatible in many cases of practical interest, as has been pointed out previously [Ref. 2, p. 2].

II. An Integral Representation

If the half toroid simulator is in the vertical position and image theory is invoked it is clear that the magnetic field at the surface of an infinite semicircular cylinder will be the same as that at the surface of an infinite circular cylinder coaxial with a current-carrying loop. Such a structure is shown schematically in figure 1. Figure 1 is presented mainly to indicate the coordinate systems we will be using.

One representation of the magnetic vector potential due to the current in the loop shown in Figure 1 is [Ref. 5]:

$$A_{\phi}^{1}(\rho,z) = \frac{\mu_{o}^{aI}}{\pi} \int_{0}^{\infty} I_{1}(k\rho)K_{1}(ka)\cos(kz)dk \quad \rho < a .$$
(1)

We must add to this the magnetic vector potential due to the currents flowing on the infinite cylinder. This latter potential may be written in the form

$$A_{\phi}^{2}(\rho,z) = \frac{\mu_{o}^{aI}}{\pi} \int_{0}^{\infty} f(k)K_{1}(k\rho)\cos(kz)dk \qquad \rho > b \qquad (2)$$

But

and, since the cylinder is assumed to be perfectly conducting, B_{ρ} is zero at its surface. This condition may be satisfied by setting

$$f(k) = -\frac{I_1(kb)K_1(ka)}{K_1(kb)} .$$
 (3)

If one substitutes equation (3) in equation (2) and then computes

$$B_{z}(\rho,z) = \frac{1}{\rho} \frac{\partial}{\partial \rho} \rho (A_{\phi}^{1} + A_{\phi}^{2}) , \qquad (4)$$

he finds

$$B_{z}(\rho,z) = \frac{\mu a I}{\pi} \int_{0}^{\infty} k \left[K_{1}(ka) I_{0}(k\rho) + \frac{K_{1}(ka) I_{1}(kb) K_{0}(k\rho)}{K_{1}(kb)} \right] \cos(kz) dk \quad \rho < a \quad (5)$$

$$B_{\rho} \propto \frac{\partial R_{\phi}}{\partial z}$$

By using the Wronskian relation for the modified Bessel functions the above expression for $B_z(\rho,z)$ at the surface of the cylinder ($\rho = b$) may be reduced to

$$B_{z}(b,z) = \frac{\mu a I}{\pi b} \int_{0}^{\infty} \frac{K_{1}(ka)}{K_{1}(kb)} \cos(kz) dk \qquad (6)$$

The current density on the cylinder is given by

$$K_{\phi}(z) = B_z/\mu_o$$

We may define a parameter

and introduce a normalized distance along the cylinder,

d ≡ z/a

to express the normalized surface current density on the cylinder

$$f(z) \equiv \frac{2aK_{\phi}(z)}{I}$$

in the form

$$f(d,\alpha) = \frac{2}{\alpha\pi} \int_{0}^{\infty} \frac{K_{1}(x)}{K_{1}(\alpha x)} \cos(xd) dx$$
(7)

As α approaches zero (i.e. as the loop recedes from the cylinder), equation (7) reduces to

$$f(d,\alpha) \xrightarrow{\alpha \to 0} \frac{2}{\pi} \int_{0}^{\infty} x K_{1}(x) \cos(xd) dx = (d^{2} + 1)^{-3/2} .$$
(8)

The central member of the above relation comes from the small argument asymptotic form of the modified Bessel function, while the equality is a special case of the equation given in Reference 6. Equation (8) is also to be expected from a consideration of the magnetic field on the axis of an isolated circular loop.

By using equations (7) and (8), one can define a "relative error" in the normalized current density, i.e. the difference between the true current density and that which would be present if the loop were extremely remote. This "relative error" will be defined by

$$\Delta(d,\alpha) \equiv (1 + d^2)^{3/2} f(d,\alpha) - 1$$
(9)

Equations (7) and (9) were evaluated numerically.

III. Numerical Calculations

As stated at the end of the previous section the numerical computations were based on equation (7). Accuracy in the numerical results is assured by the fact that the integration intervals between 0 and $5/\alpha$ were kept small compared to 1/z and the contribution to the integral from the region $x > 5/\alpha$ was approximated, using the large argument asymptotic forms of the modified Bessel functions, by

$$\int_{5/\alpha}^{\infty} \frac{K_1(x)}{K_1(\alpha x)} \cos(xd) dx \approx \sqrt{\alpha} e^{-(1-\alpha)^{5/\alpha}} \cdot \frac{(1-\alpha)\cos(5d/\alpha) - d \sin(5d/\alpha)}{(1-\alpha)^2 + d^2}$$

In Table 1 are the values of $f(d,\alpha)$ for various values of d and α . The corresponding values of $\Delta(d,\alpha)$ are given in Table 2. Figure 2 contains plots of $f(d,\alpha)$ as a function of d with α as a parameter while Figure 3 contains the same information plotted as a function of α with d as a parameter. Figures 4 and 5 are plots of the $\Delta(d,\alpha)$ corresponding to Figures 2 and 3.

α d	0	• 2.	.4	.6	.8	1.0	1.2
• 0	1.000	.943	.800	.631	.476	.354	.262
.02	1.002	.945	.801	.630	.476	.353	.262
.04	1.007	.948	.803	.630	.475	.352	.261
.06	1.014	.953	.805	.630	.473	350	.259
.08	1.022	960	.807	.630	472	348	257
10	1 032	967	811	629	470	346	255
.12	1.043	.976	.814	.629	.467	343	.253
.14	1.056	.985	.818	.629	.465	.340	.250
.16	1.069	.995	.822	.628	.462	.337	.248
.18	1.083	1.007	.827	.628	.460	334	. 245
.20	1.099	1.019	.831	.627	.457	.331	. 241
.20	1.116	1.032	.836	.626	.453	327	.238
.24	1.134	1.045	.841	.625	.450	323	.234
.26	1,154	1,060	.846	.623	.446	.318	.231
.28	1,175	1.076	.852	.622	.441	.314	.227
.30	1.198	1.092	.857	.620	.437	.309	.223
.32	1.223	1.110	.862	.617	.432	.304	.218
.34	1.249	1.128	.867	.615	.427	.299	.214
.36	1.277	1.148	.872	.612	.421	.294	.209
•38 ·	1.308	1.168	.877	.608	.415	.288	.205
. 40	1.341	1.190	.881	.603	.409	.282	.200
• .42	1.375	1.212	.885	.599	.402	.276	.195
.44	1,414	1.236	.889	.593	.395	.269	.189
.46	1.455	1.261	.892	.587	.387	.262	.184
.48	1.499	1.287	.894	.580	.378	.255	.178
.50	1.547	1.315	. 895	.572	.370	.248	.173
.52	1.600	1.343	.896	.563	.360	.240	.167
.54	1.657	1.373	.895	.553	.350	.232	.161
.56	1.720	1.403	.892	.543	.340	.224	.155
.58	1.788	1.435	.888	.531	.329	.215	.149
.60	1.864	1.467	.882	.518	.318	.207	.142
.62	1.948	1.500	.874	.503	.306	.198	.136
.64	2.041	1.534	.863	.488	.294	.188	.129
.66	2.147	1.567	.850	.471	.281	.179	.122
.68	2.264	1.599	.833	.453	.268	.169	.115
.70	2.398	1.629	.813	.433	.254	.160	.107
.72	2.551	1.657	.789	.412	.240	.150	.100
.74	2.727	1.680	.761	.389	.225	.140	.093
.76	2.933	1.696	.728	.365	.210	.130	.086
.78	3.176	1.703	.691	.339	.194	.120	.078
.80	3.468	1.696	.648	.312	.178	.110	.071
.82	3.824	1.673	.601	.284	.161	.099	.064
.84	4.268	1.628	.549	.255	.144	.089	.057
.86	4.840	1.555	.492	.225	.126	.078	.050
.88	5.601	1.447	.430	.193	.109	.067	.043
.90	6.665	1.301	.364	.162	.091	.056	.036
.92	8.260	1.11	•294	.129	.073	.045	.029
•94	10.916	.8/9	.222	.097	.054	.034	.022
•96	10.225	.608	.149	.064	.036	.023	.015
. 98	32.145	.310	.074	.032	.018	011	. 007

Table 1: f(d, a); normalized surface current on an infinite cylinder

Table 2: $\Delta(d, \alpha)$; "Relative Error" in normalized surface current

αd	0	.2	• 4	.6	.8	1.0	1.2
0	.000	.000	.000	.000	.000	.000	.000
.02	.002	.002	.001	000	001	002	002
.04	.007	.006	.003	001	003	005	006
.06	.014	.011	.005	001	006	010	012
.08	.022	.018	.009	001	010	015	019
.10	.032	.026	.013	002	014	022	027
.12	.043	.035	.017	002	018	029	036
.14	.056	.045	.022	003	023	037	046
.16	.069	.056	.027	004	029	046	056
.18	.083	.068	.033	005	035	055	068
.20	.099	.080	.039	006	041	065	080
.22	.116	.094	.045	007	048	076	093
.24	.134	.109	.051	009	056	087	107
.26	.154	.124	.057	011	064	099	121
.28	.175	.141	.064	014	073	112	136
.30	.198	.158	.070	017	082	125	152
.32	.223	.177	.077	021	093	139	168
.34	.249	.196	.083	025	104	154	185
.36	.277	.217	.089	030	115	170	202
.38	.308	.239	.095	036	128	186	220
.40	.341	.262	.101	043	141	202	239
. 42-	: .376	.286	106	051	156	220	258
44	- 414	.311	.111	059	171	239	278
.46	.455	.338	.114	069	188	258	299
.48	.499	. 365	.117	081	205	278	320
.50	.547	.394	.119	093	224	299	341
.52	.600	.424	.119	107	243	320	363
.54	.657	.456	.118	122	264	343	386
.56	.720	.488	.115	139	285	366	409
.58	. 788	.522	.110	158	308	391	433
.60	.864	.556	.102	179	332	415	- .458
.62	.948	.591	.092	201	357	441	483
.64	1.042	.627	.079	226	382	467	509
.66	1.147	.662	.062	253	409	494	536
•68	1.264	.696	.041	282	437	521	563
.70	1.398	.728	.015	314	466	548	591
•72	1.551	.757	015	347	496	576	619
•74	1.727	.781	050	383	527	604	646
.76	1.933	.798	090	422	560	632	6/4
.78	2.176	.806	13/	462	593	661	/01
.80	2.468	.799	190	505	627	690	729
.82	2,824	.//5	249	549	002	/19	/56
.84	3.268	./26	314	390			/83
•86	3.840	.649	300	044	/33	//9	810
.88	4,601	.535	403 E/E	093	- 200	- 01U	-,03/
.90	5,665	.380	343	/44	009	041	003
.92	/.260	.1/9		- 01.6	040	0/2	090
.94	9.910 15 005	U00 255	- 01/	040 000	000	- 024	· - 91/
• 96	13.223	300	- 009	- 070	- 949	- 930	744
• 78	JL.14J	0/1			702	200	714



Figure 1: An infinite cylinder within a circular loop.







Figure 3. $f(d,\alpha)$ versus α with d as a parameter.



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